Conformational Stability of Ibuprofen: Assessed by DFT Calculations and Optical Vibrational Spectroscopy

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ABSTRACT: A thorough conformational analysis of ibuprofen [2-(4-isobutylphenyl) propionic acid] was carried by out, using density functional theory (DFT) calculations coupled to optical vibrational spectroscopy (both Raman and FTIR). Eight different geometries were found to be energy minima. The relative orientations of the substituent groups in the ibuprofen molecule, which can be considered as a para-substituted phenyl ring, were verified to hardly affect its conformational stability. The internal rotations converting the calculated conformers of ibuprofen were studied and the intramolecular interactions governing the conformational preferences of the molecule were analyzed by quantitative potential energy deconvolution using Fourier type profiles. The harmonic vibrational frequencies and corresponding intensities were calculated for all the conformers obtained, leading to the assignment of the spectra, and evidencing the sole presence of one of the lowest energy conformers in the solid state. Vibrational spectroscopic proof of intermolecular hydrogen bonds between the carboxylic groups of adjacent ibuprofen molecules, leading to the formation of dimers, was also obtained. © 2007 Wiley-Liss, Inc. and the American Pharmacists Association J Pharm Sci 97:845-859, 2008

Keywords: ibuprofen; DFT calculations; Raman spectroscopy; FTIR spectroscopy; conformational analysis; rotational isomerism

INTRODUCTION

The knowledge of the conformational preferences of some nonsteroidal anti-inflammatory drugs (NSAID's), particularly the substituted α -arylpropionic acids, is of the utmost importance for a better understanding of the structure–activity relationships (SAR's) underlying their biological activity, as well as of their mechanism of action,

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which can be related to the ability to inhibit prostaglandin synthesis.^{1–3} In fact, several studies are reported which establish the monocarboxylic acid group responsible for the antiinflammatory, analgesic, and antipyretic properties of these compounds.^{4–6} However, different ring substitution is known to give rise to distinct pharmacological properties, probably due to pharmacokinetic and pharmacodynamic variations. Moreover, the conformational behavior of the drug determines the chemical and/ or physical mechanisms (i.e., intermolecular interactions) controlling its release into the body from a particular delivery system, and consequently its bioavailability.

Different physico-chemical methods have been employed for obtaining structural information on pharmacologically relevant systems, as well as for understanding the interactions between the therapeutically active agent and the excipients and/or carriers. Among these, Raman spectroscopy has lately become more and more used as a fast and nondestructive technique for the study of this kind of compounds, since it allows a direct observation of the sample, without any special preparation procedures that could interfere with its conformational characteristics.^{7–9} Moreover, the information yielded by Raman spectroscopy is complementary to the one gathered by infrared analysis, both methods providing a complete vibrational information on a particular molecule.

Ibuprofen (Fig. 1) (2-(4-isobutylphenyl) propionic acid), the parent compound of the "profen"- type NSAID's, is a weakly acidic ($pK_a = 4.6$), poorly water-soluble drug (water solubility ≈ 0.05 mg/mL at 25°C),¹⁰ frequently used for the treatment of painful and inflammatory conditions such as rheumatoid arthritis, osteoarthritis, and ankylosing spondvlitis.¹¹ This compound comprises a chiral centre at the α -carbon and can therefore exist as R(-) or S(+) enantiometric forms. The inactive R(-)enantiomer undergoes a unidirectional chiral inversion *in vivo* to S(+), which is the species displaying anti-inflammatory activity.¹²⁻¹⁵ Following a single dose administration of immediate-release ibuprofen preparations, the peak plasma drug concentration was observed at 3 h postdose. Moreover, this drug is characterized by a rapid onset of pharmacological actions, with relatively little associated gastric disturbance, the blood plasma half-life lying between 2 and 3 $h^{16,17}$ This short half-life,

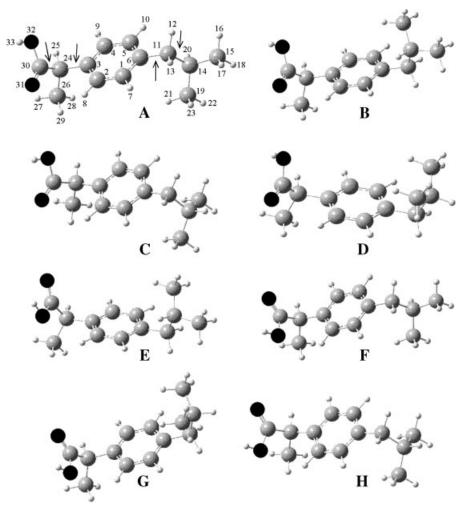


Figure 1. (A–H) Schematic representation of the eight most stable conformers of the ibuprofen molecule.

coupled to the low single administration dosage necessary, renders ibuprofen a good candidate for the development of new controlled-delivery formulations. $^{18-21}$

The development of distinct ibuprofen hydrophilic matrix tablets using different excipients (e.g., swellable cellulose polymers) is presently the object of vigorous research, as extended-release (either constant or pulsed) dosage forms of a particular drug may often be beneficial. Thus, the knowledge of intermolecular interactions between ibuprofen and these polymers/excipients, which modulates the *in vivo* drug release process, is of the utmost importance. This may be achieved through optical vibrational spectroscopy (both Raman and FTIR), once the conformational behavior of the pure drug in the solid state is known. The present study aims at achieving this goal, which will hopefully allow carrying out future studies on ibuprofen tablets (composed of distinct drug/polymer/excipient mixtures).

To the best of our knowledge, the only conformational study of ibuprofen reported to date was done by crystallographic database searching and potential energy calculations using the semi-empirical AM1 method.²²

In the present work, a complete conformational study of ibuprofen was thus undertaken by quantum mechanics density functional theory (DFT) calculations coupled to Raman spectroscopy. The FTIR spectrum of pure ibuprofen was also analyzed, in order to explore the well known complementary between these two optical vibrational spectroscopy techniques.

MATERIALS AND METHODS

Material

Ibuprofen batch no. 9907257 was purchased from Knoll, Nottingham, England.

Methods

Raman Spectroscopy

The Raman spectra were obtained on a triple monochromator Jobin-Yvon T64000 Raman system (focal distance 0.640 m, aperture f/7.5) equipped with holographic gratings of 1800 grooves \cdot mm⁻¹. The premonochromator stage was used in the subtractive mode. The detection system was liquid nitrogen cooled nonintensified

 578×385 pixel (1/2") charge coupled device (CCD) chip. A coherent (model Innova 300-05) Ar⁺ laser was used as light source, the output of which at 514.5 nm was adjusted to provide 35 mW at the sample position. A 90° geometry, between the incident radiation and the collecting system, was employed. The entrance slit was set to 200 µm and the slit between the premonochromator and the spectrograph was opened to 12 mm. An integration time of 3 s and 10–15 scans were used in all experiments.

Samples were sealed in Kimax glass capillary tubes of 0.8 mm inner diameter. Under the abovementioned conditions, the error in wave numbers was estimated to be within 1 cm^{-1} .

FTIR Spectroscopy

Infrared spectra of ibuprofen in KBr disks (ca. 5% (w/w)) were recorded at room temperature on a Nicolet Model 740 FTIR spectrometer, in the range 400–4000 cm⁻¹, using a globar source, a Ge/KBr beamsplitter, a DTGS detector. The spectra were collected in 32 scans to a 16384 data points file (resolution ca. 2 cm⁻¹) and subject to a Happ-Genzel apodization. The errors in wave numbers were estimated to less than 1 cm⁻¹.

DFT Calculations

The molecular orbital calculations were carried out with the GAUSSIAN 03W program,²³ within the DFT approach, using the B3LYP method, which includes a mixture of Hartree-Fock (HF) and DFT exchange terms. The gradient-corrected correlation functional was used,^{24,25} parameterized after Becke,^{26,27} along with the double-zeta split valence basis set $6-31G^*$.²⁸

Molecular geometries were fully optimized by the Berny algorithm, using redundant internal coordinates²⁹: the bond lengths to within ca. 0.1 pm and the bond angles to within ca. 0.1°. The final root-mean-square (rms) gradients were always less than 3×10^{-4} Hartree · bohr⁻¹ or Hartree · radian⁻¹. In order to study the barriers to internal rotation, the geometries were optimized for different fixed internal rotation angles.

The potential-energy profiles for rotation around several bonds within the molecule were obtained by scanning the correspondent dihedral angles by 15° steps. The quantitative deconvolution of these profiles was based on least-squares fitted Fourier type functions of a relevant torsional angle,

$$V = V_0 + \sum_{n=1}^4 \frac{1}{2} V_n [1 - \cos(n\tau)] + \sum_{m=1}^4 V'_m \sin(m\tau)$$

where τ represents the $O^{31}C^{30}C^{24}C^3$, $HC^{24}C^3$ C^4 , $C^1C^6C^{11}C^{14}$, or $C^6C^{11}C^{14}H$ dihedrals (Fig. 1), and V are functional values that correspond to potential energy differences relative to a reference value $(V_0, \text{ the energy corresponding to a dihedral})$ angle of 0°). According to the specific profile under study, different combinations of cosine and sine term were used (from 3 to 6).

Orbital interactions were determined using the natural bond orbital (NBO) approach from the donor-acceptor viewpoint³⁰ applied to the wave functions, at the B3LYP/6-31G* level of calculation. Routines for this kind of calculation are included in the GAUSSIAN 03W package, which convert the DFT molecular orbitals in a set of NBO orbitals, which constitute a hypothetical Lewis structure with strictly localized electron pairs. In this NBO formulation, delocalization arises from interactions between occupied bonding and antibonding orbitals and is represented by off-diagonal terms in the Kohn-Sham matrix.³¹

RESULTS AND DISCUSSION

Conformational Analysis

Ibuprofen can adopt different conformations, mainly by varying the dihedral angles around the $C^{30}-C^{24}$, $C^{24}-C^{3}$, $C^{6}-C^{11}$, and $C^{11}-C^{14}$ bonds (Fig. 1A). Moreover, rotational isomerism is also possible within the O=C–O–H group, giving rise to either s-cis (0°) or s-trans (180°) stable geometries. Several studies performed on carboxylic containing molecules³²⁻³⁴ have demonstrated that, in the absence of intramolecular stabilizing interactions, the s-cis arrangements were found to be significantly more stable (ca. 20 kJmol⁻¹) than their s-*trans* counterparts. Thus, in the present study only those geometries displaying O=C-O-H dihedrals equal to ca. 0° were considered.

The eight different optimized conformations represented in Figure 1 were found to correspond to minima in the potential energy surface. This was verified by the absence of DFT calculated imaginary (negative) frequencies. Table 1 comprises the conformational energy differences. dipole moments, rotational constants, and thermochemical data (at 298.15 K and 1 atm) for all the ibuprofen conformers presently calculated.

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$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	Conformer ^a	$0 = C^{30}C^{24}C^{3}$	$\mathrm{HC}^{24}\mathrm{C}^{3}\mathrm{C}^{4}$	$C^{1}C^{6}C^{11}C^{14}$	$\mathrm{C}^{6}\mathrm{C}^{11}\mathrm{C}^{14}\mathrm{H}$		Population b (%)	ALZPVE (kJmol ⁻¹)		1 do (kJmol ⁻¹)	من (kJmol ⁻¹)	Population d (%)	$\mu^{e} \; (\mathrm{D})$	A; B; C (GHz)
$ \begin{array}{llllllllllllllllllllllllllllllllllll$	Α	-89.6	-10.2	105.1	54.9	0.00	37.7	0.00	0.00	00.00	00.00	41.9	1.557	1.569; 0.240; 0.235
			171.0	-74.5										
$ \begin{array}{l lllllllllllllllllllllllllllllllllll$	В	-89.5	171.5	105.6	55.0	0.05	36.9	-0.08	-0.07	-0.05	-0.02	42.2	1.401	1.326; 0.253; 0.245
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$			-10.2	-73.8										
$ \begin{array}{c ccccccccccccccccccccccccccccccccccc$	C	-89.1	-9.5	89.9		4.12	7.2	4.93	4.54	-2.22	6.75	2.7	1.489	1.396; 0.261; 0.258
$ \begin{array}{cccccccccccccccccccccccccccccccccccc$	D	-89.2	-11.1	-90.3		4.16	7.0	4.89	4.50	-2.17	6.68	2.8	1.502	1.191; 0.280; 0.270
	Е	99.2	-1.3	-75.7		5.16	4.7	5.12	5.07	-0.48	5.55	4.5	1.782	1.327; 0.252; 0.245
	F	98.1	-3.6	103.2		5.21	4.6	5.03	5.00	-0.19	5.19	5.2	1.781	1.589; 0.238; 0.234
	Ċ	97.8	-4.2	-90.2	180.0	9.19	0.9	9.90	9.43	-2.52	11.95	0.3	1.752	1.187; 0.278; 0.270
^a See Figure 1. ^b According to ΔE values. ^c $\Lambda E_{correct}$ zero-noint vibrational energy corrected relative energies	Н	98.1	-3.3	90.2		9.25	0.9	9.99	9.59	-2.15	11.74	0.4	1.834	1.416; 0.260; 0.255
b According to ΔE values. $^cAE_{max}$ zero-mint vibrational energy connected relative energies	^a See Fig	ture 1.												
	$^{o}According ^{o}According $	ing to <i>ΔE</i> val. . zero-point v	ues. ibrational e	merev correc	ted relative	energies.								

^{*d*}According to ΔG values. ^{*e*}I D = 1/3 × 10⁻² C · m.

Coordinate ^a	Experimental	$Calculated^b$	Calculated
Pond longth (nm)			
Bond length (pm)			
$C^{30} - O^{32}$	1.306	1.356	1.355
$C^{30}-O^{31}$	1.204	1.211	1.213
C^{30} – C^{24}	1.503	1.526	1.523
C^{24} – C^{26}	1.500	1.535	1.538
$C^{24}-C^{3}$	1.525	1.530	1.528
C^3-C^4	1.374	1.401	1.400
C^4 – C^5	1.376	1.393	1.394
C^5 – C^6	1.392	1.402	1.402
$C^{6}-C^{1}$	1.380	1.400	1.401
C^1 – C^2	1.396	1.395	1.395
$C^{2}-C^{3}$	1.380	1.400	1.401
$C^{6}-C^{11}$	1.493	1.514	1.514
C^{11} – C^{14}	1.529	1.550	1.550
C^{14} – C^{15}	1.508	1.535	1.535
C^{14} – C^{19}	1.519	1.534	1.535
O^{32} – H^{33}	0.963	0.976	0.976
C^{24} – H^{25}	1.091	1.096	1.095
C^{26} – H^{29}	1.081	1.093	1.093
$C^{26}-H^{27}$	1.053	1.095	1.096
$C^{26}-H^{28}$	1.073	1.094	1.094
C^4 – H^9	1.103	1.087	1.088
$C^{5}-H^{10}$	1.041	1.088	1.088
C^1-H^7	1.065	1.088	1.088
$\tilde{C}^2 - H^8$	1.077	1.086	1.087
$C^{11}-H^{13}$	1.101	1.099	1.099
$C^{11}-H^{12}$	1.102	1.099	1.099
$C^{14}-H^{20}$	1.085	1.100	1.100
$C^{15}-H^{17}$	1.061	1.098	1.098
$C^{15}-H^{16}$	1.062	1.097	1.097
$C^{15}-H^{18}$	1.097	1.096	1.096
$C^{19}-H^{21}$	1.067	1.095	1.095
$C^{19}-H^{22}$	1.007	1.096	1.096
$C^{19}-H^{23}$	1.033	1.098	1.098
0 11	1.011	1.050	1.050
Bond angle (°)			
O^{32} - C^{30} - O^{31}	123.4	122.5	122.4
O^{32} - C^{30} - C^{24}	115.4	111.7	111.9
O^{31} - C^{30} - C^{24}	121.1	125.7	125.7
C_{30} – C_{24} – C_{26}	111.7	110.2	110.4
C_{30} - C_{24} - C_{3}	106.7	108.5	109.4
C_{26} - C_{24} - C_{3}	114.4	114.2	112.3
C_{24} - C_{3} - C_{4}	120.9	120.1	120.5
$C^{3}-C^{4}-C^{5}$	121.6	121.0	120.8
$C^{4}-C^{5}-C^{6}$	120.7	121.2	121.2
$C^{5}-C^{6}-C^{1}$	118.0	117.6	117.7
$C^{6}-C^{1}-C^{2}$	120.7	121.4	121.4
U -U -U			
C = C = C $C^{1} = C^{2} = C^{3}$	120.7	120.7	120.6

Table 2.	Experimental (Single-Crystal Pulsed
Neutron D	iffraction ³⁵) and Calculated (B3LYP/6-31G*)
Geometric	al Parameters for Ibuprofen

 Table 2. (Continued)

$\operatorname{Coordinate}^{a}$	Experimental	$Calculated^{b}$	$Calculated^{c}$
C^2 - C^3 - C^{24}	120.9	121.7	121.1
$C^{5}-C^{6}-C^{11}$	120.2	120.7	120.7
$C^{1}-C^{6}-C^{11}$	121.8	121.7	121.6
C^6 - C^{11} - C^{14}	113.9	114.6	114.6
C^{11} - C^{14} - C^{15}	110.1	110.3	110.3
C^{11} - C^{14} - C^{19}	111.5	112.0	112.0
C^{15} - C^{14} - C^{19}	111.5	111.0	111.0
Torsional angle (°)			
O^{31} - C^{30} - C^{24} - C^{3}	-89.6	-95.5	-89.6
O^{32} - C^{30} - C^{24} - C^{3}	88.7	83.0	89.0
$H^{33}-O^{32}-C^{30}-O^{31}$	-3.3	1.2	1.5
$H^{33}-O^{32}-C^{30}-C^{24}$	-175.1	-176.6	-177.2
O^{31} - C^{30} - C^{24} - C^{26}	36.0	30.2	34.4
H^{25} - C^{24} - C^{3} - C^{4}	19.8	19.8	-10.2
C^{26} - C^{24} - C^{3} - C^{4}	140.5	140.6	110.6
C^{30} - C^{24} - C^{3} - C^{4}	-95.5	-96.0	-126.6
C^{24} - C^{3} - C^{4} - C^{5}	-177.7	178.5	-178.2
C^1 C^6 C^{11} C^{14}	102.1	103.9	105.1
$C^6-C^{11}-C^{14}-C^{15}$	168.5	171.7	172.4
C^6 - C^{11} - C^{14} - H^{20}	50.4	54.3	54.9
C^6 - C^{11} - C^{14} - C^{19}	-67.1	-64.0	-63.4
$\Delta E \; (\mathrm{kJmol}^{-1})^d$	5.12	2.97	0.00

^aSee Figure 1 for atom numbering.

^bComplete optimization based on the experimental geometry, for a fixed $H^{25}-C^{24}-C^{3}-C^{4}$ dihedral at 19.8°.

^cMost stable calculated conformer (A). ^dCalculated (B3LYP/6-31G^{*}) relative energy values.

It is noteworthy that the two sets of dihedrals $HC^{24}C^{3}C^{4} = -10.2^{\circ}$ or 105.1° and $C^{1}C^{6}C^{11}$ $C^{14} = 171.0^{\circ}$ or -74.5° correspond to conformer A, since they yield identical structures, the same occurring for B (Tab. 1).

The ibuprofen molecule can be considered as a para-substituted phenyl ring. Interestingly, the relative orientation of the substituents hardly affects the conformational stability of this system. In fact, these can be either below or above the ring plane, or even in opposite sides (Fig. 1). Actually, the rotations around the C^{24} – C^3 and C^6 – C^{11} bonds are not correlated, which is evidenced by comparing conformers (both geometries and energies) A versus B, C versus D, E versus F, and G versus H (Fig. 1 and Table 1). Consequently, the existence of not more than four energetically distinct conformations may be considered: AB ($\Delta E \approx$ 0 kJmol⁻¹), CD ($\Delta E \approx 4.1$ kJmol⁻¹), EF ($\Delta E \approx$ 5.2 kJmol⁻¹), and GH ($\Delta E \approx 9.2$ kJmol⁻¹) (Tab. 1). These energy differences correspond to room temperature populations of 75, 14, 9, and 2%, respectively.

However, although the two structures in each of these groups are virtually degenerated, their rotational constants and, in some cases, their dipole moments do not follow the same energetic trend (Tab. 1). This would be expected, since the change in the relative orientation of the substituents, which leads to different rotational constants and dipole moments, was found not to affect the corresponding conformational stability (Fig. 1).

The estimation of the ΔE based populations assumed similar entropy contributions to the free energy of the distinct conformers. Nevertheless, this assumption is not strictly correct since the smaller conformational flexibility of C, D, G, and H ($C^6C^{11}C^{14}$ H ca. 180°) provides less significant entropy terms $(T\Delta S)$ and, accordingly, an increase of the corresponding free energy differences relative to the most stable conformers: AB $(\Delta G \approx 0 \text{ kJmol}^{-1})$, CD $(\Delta G \approx 6.7 \text{ kJmol}^{-1})$, EF $(\Delta G \approx 5.4 \text{ kJmol}^{-1})$, and GH $(\Delta G \approx 11.8 \text{ kJmol}^{-1})$ (Tab. 1). These free energy differences correspond to room temperature Gibbs populations of 84, 5, 10, and 1%, respectively. Moreover, if entropy contributions are considered, the B conformer becomes the most stable one $(\Delta G_{\rm A-B} \approx$ -0.02 kJmol^{-1}).

In aqueous solutions, however, these populations are expected to be slightly altered, since the conformers displaying higher dipole moments (E, F, G, and H) are energetically favored. The effect of the aqueous environment can even overrule the energetic degeneracy discussed above: particularly for the AB pair, the population of A $(\mu = 1.557 \text{ D})$ will certainly increase relative to B $(\mu = 1.401 \text{ D})$.

Table 2 comprises the single-crystal pulsed neutron diffraction refined geometry reported for ibuprofen³⁵, as well as some B3LYP/6-31G* optimized geometries (results for the other conformers and for angle parameters involving hydrogen atoms are available from the authors upon request). The DFT structural parameters presently obtained are in guite good agreement with the experimental ones, although there are some discrepancies both in bond distances (e.g., C^{30} – O^{32} and C^{24} – C^{26}) and angles (e.g., O^{31} – $C^{30}-C^{24}$) which can be a consequence of the crystal packing. The energy difference between the most stable ibuprofen conformer (B3LYP/6-31G* total geometry optimization) and the one obtained from the experimental geometry (B3LYP/6-31G* single-point calculation) is 5.12 kJmol⁻¹, due to two main factors: (i) conformational rearrangements around the $C^{24}-C^3$ bond (ca. 30° internal rotation, Table 2) which are required for an efficient molecular packing in the crystalline structure; (ii) the optimization process did not allow for the relaxation of the other geometrical parameters. In order to abolish the latter and obtain a more accurate value for the energy gap associated to that rearrangement, a total geometry optimization was performed for a fixed $H^{25}-C^{24}-C^3-C^4$ dihedral angle (Tab. 2), which yielded energy of 2.97 kJmol⁻¹.

Rotational isomerism in this type of compounds, comprising aromatic, carboxylic and alkylic groups is dependent on several factors, namely steric, dipolar, mesomeric, and hiperconjugative effects, and hydrogen bond interactions. Furthermore, the relative importance of intra versus intermolecular interactions (e.g., dimer formation) has often proved, in several systems, to be determinant of their conformational preferences, either as pure compounds or in solution.

Potential energy profiles for internal rotation around different bonds within the ibuprofen molecule ($C^{30}-C^{24}$, $C^{24}-C^{3}$, $C^{6}-C^{11}$, and $C^{11}-C^{14}$ bonds) were obtained, by scanning the corresponding torsional angles.

Rotation around the $C^{30}-C^{24}$

The C^{30} - C^{24} rotation converts conformers A to F (Fig. 2a). The energy difference between these two arrangements $(\Delta E_{\rm F-A})$ and the corresponding clockwise internal rotation barrier $(A \rightarrow F)$ being 5.21 kJmol⁻¹ and 9.9 kJmol⁻¹, respectively. The anticlockwise rotation, in turn, has a 11.3 kJmol⁻¹ barrier. The presence of the aromatic ring is responsible for obvious differences in the conformational behavior of the ibuprofen propionic moiety, when compared to previous studies on both propionic and 2-methylpropionic acids.^{36–38} Actually, latter the preferred conformations around the C α -C bond are those displaying α substituents with either syn or skew orientations relative to the C=O bond (CCC=O equal to 0° or $\pm 120^{\circ}$, respectively). The analogous geometries for ibuprofen were also found to be favored, as the Fourier term in V_3 (Fig. 2b) displays an intermediate value $(2.12 \text{ kJmol}^{-1})$, despite the largely dominant contribution being represented by a cosine component in V_2 (-5.66 kJmol⁻¹). This term exhibits maxima for C³C²⁴C³⁰=O³¹ equal to 0° and 180° , and minima for -90° and 90° , reflecting the preference for geometries that favor π -delocalization between the OC=O group and the

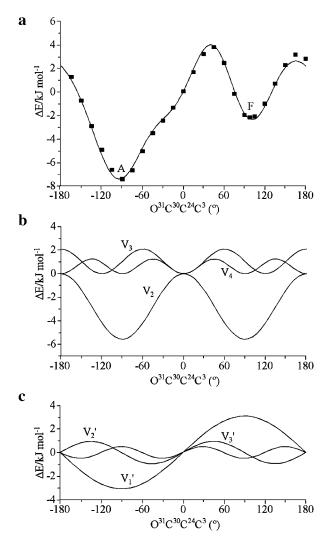


Figure 2. Optimized (B3LYP/6-31G^{*}) conformational energy profile for the internal rotation around the $O^{30}-C^{24}$ bond of ibuprofen (a) and its Fourier deconvolution (b and c). $V_2 = -5.66 \text{ kJmol}^{-1}$, $V_3 = 2.12 \text{ kJmol}^{-1}$, $V_4 = 1.27 \text{ kJmol}^{-1}$, $V_1' = 3.26 \text{ kJmol}^{-1}$, $V_2' = 0.87 \text{ kJmol}^{-1}$, and $V^3 = 0.53 \text{ kJmol}^{-1}$.

aromatic ring, which are face-to-face (in-phase) in these conformations (Fig. 2b) (as opposed to the tilted, out-of-phase, relative orientations occurring in the other geometries). In fact, the results yielded by the NBO approach and donor–acceptor analysis based on localized orbitals,³⁰ allow to conclude that $\pi(C^2-C^3) \rightarrow \pi^*(C^{30}=O^{31})$ is the key interaction occurring in these conformations, with a calculated stabilization energy of 8.54 kJmol⁻¹. Moreover, the $\pi(C^{30}=O^{31}) \rightarrow \sigma^*(C^3-C^{24})$ and $\sigma(C^3-C^{24}) \rightarrow \pi^*(C^{30}=O^{31})$ interactions also contribute to the described stabilization. Thus, two minima were obtained for the ibuprofen $C^{30}-C^{24}$

rotational profile, for $C^3C^{24}C^{30}{=}O^{31}$ equal to -89.6° and $98.1^\circ.$

Furthermore, these two minima are not equivalent ($\Delta E_{\rm F-A} = 5.21 \text{ kJmol}^{-1}$). In fact, the V_1 sine term (3.23 kJmol⁻¹, Fig. 2c) favors the conformations displaying a negative $C^3C^{24}C^{30}=O^{31}$ dihedral angle, that is, for an equal orientation of the carbonyl and the α -methyl group. Actually, these geometries allow the formation of stabilizing C= O^{-1} H close contacts, either with H²⁷, H²⁹, or H⁸ (Fig. 3). In fact, H₈ displays the highest natural atomic charge (0.247) as compared to the other aromatic hydrogens (ca. 0.234) (Fig. 3).

Rotation around C^{24} – C^3

The rotation C^{24} – C^3 interconverts conformers A and B (Fig. 4a), $HC^{24}C^3C^4$ being equal to -10.2° and 171.5° , respectively. The corresponding internal conversion barrier is 13.5 kJmol^{-1} , both conformers having an equivalent potential energy ($\Delta E_{B-A} = 0.05 \text{ kJmol}^{-1}$). This evidences that the conformational stability of the ibuprofen molecule is not affected by the relative orientation of the substituents—either below or above the ring plane (conformer B, Fig. 1) or to opposite sides of this plane (conformer A, Fig. 1).

In the light of this result, ibuprofen can be regarded as both an α,α -disubstituted toluene on C²⁴, and an α -monosubstituted toluene on C¹¹. Although hiperconjugative interactions constitute an acceptable basis for explaining the barriers for internal rotation around the C(sp²)–C(sp³) bond in α -substituted toluenes, this effect alone does not account satisfactorily for the conformational behavior of the more heavily substituted toluenes. In particular, when either heteroatoms or bulky α -substituents are present, other effects, such as steric and electrostatic interactions, must be considered.³⁹

The values of the Fourier components of the potential energy profile for the $HC^{24}C^{3}C^{4}$ dihedral

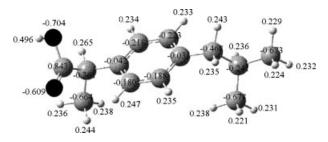


Figure 3. Calculated (B3LYP/6-31G^{*}) natural atomic charges for the A conformer of ibuprofen.

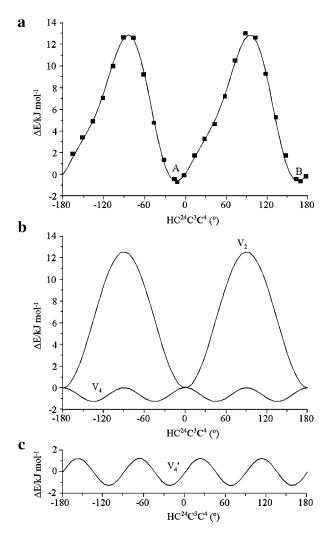


Figure 4. Optimized (B3LYP/6-31G*) conformational energy profile for the internal rotation around the $C^{24}-C^3$ bond of ibuprofen (a) and its Fourier deconvolution (b and c). $V_1 = 6.31 \text{ kJmol}^{-1}$, $V_2 = 6.31 \text{ kJmol}^{-1}$, $V_4 = -1.27 \text{ kJmol}^{-1}$, and $V'_4 = 1.24 \text{ kJmol}^{-1}$.

(Figs. 4b and c), point to a leading contribution of a cosine term in V_2 (12.62 kJmol⁻¹, Fig. 4b), reflecting a preference for *trans* and *cis* arrangements around the $HC^{24}-C^3C^4$ bond, since these conformations tend to: (i) minimise steric hindrance between the bulkiest substituent groups (α -methyl and carboxyl) and the phenyl ring, by placing the former above and below the ring; (ii) maximize the attractive hydrogen bond type interactions between the C=O carbonyl group and either H⁸ (in conformer A) or H⁹ (in conformer B), giving rise to a stabilizing six-membered intramolecular ring; (iii) favor both charge transfer $\sigma(C^{24}-H^{25}) \rightarrow \sigma^*(C^2-C^3)$ or $\sigma(C^{24}-H^{25}) \rightarrow \sigma^*(C^3-C^4)$ interactions, for $HC^{24}C^3C^4$

equal to 0° or 180° , respectively; (iv) maximize interactions between anti-bonding π orbitals from the carbonyl and phenyl groups, $\pi^*(C^{30}=O^{31})\rightarrow\pi^*(C^2-C^3).$

Also, the sine term V'_4 (1.24 kJmol⁻¹, Fig. 4c) is due to electrostatic factors occurring within the molecule. Indeed, it displays energy maxima when destabilizing interactions occur between the positively charged methyl H²⁸ or H²⁹ atoms and the aromatic H⁸ or H⁹. In contrast, it exhibits energy minima whenever the methyl group is situated in such way as to minimize this type of repulsive close contacts: (i) either above or below the ring plane, in a perpendicular relative position; (ii) or displaying C²⁶ and H²⁷ in the aromatic plane, and H²⁸ and H²⁹ symmetrically located relative to H⁸ or H⁹.

The Fourier term in V_4 (-1.27 kJ mol⁻¹, Fig. 4b), in turn, accounts for hiperconjugative stabilizing $\pi(\text{ring}) \rightarrow \sigma^*(C^{24}-C^{26})$ and $\sigma(C^{24}-C^{26}) \rightarrow \pi^*(\text{ring})$ interactions, for $HC^{24}C^3C^4$ equal to $\pm 45^\circ$ and $\pm 135^\circ$. When compared to V_2 and V'_4 , this is a much smaller contribution to the overall potential energy profile of the $-CH(CH_3)COOH$ rotor, which is indicative of the lesser significance of the hyperconjugative effect.

Rotation around $C^6 - C^{11}$

Conformers A–B (Fig. 5a), for $C^1C^6C^{11}C^{14}$ equal to -74.5° and 105.6° , respectively, are also interconverted by rotation about C^6-C^{11} , with an internal barrier of 15.1 kJmol⁻¹, which is slightly higher than the one obtained for the $C^{24}-C^3$ rotation.

Considering the values obtained for the Fourier components of the potential energy profile of this internal rotation process (Figs 5b and 5c), it can be concluded that the dominant contributions are both V_2 (-9.39 kJmol⁻¹) and V'_2 (5.27 kJmol⁻¹). The former accounts for: (i) steric hindrance minimization, when the bulky isobutyl group is away from the phenyl ring plan, corresponding to a $C^{1}C^{6}C^{11}C^{14}$ dihedral equal to $\pm 90^{\circ}$; (ii) max- $\begin{array}{ll} \mbox{imization of the } \sigma(\mbox{ring}) \rightarrow \sigma^*(\mbox{C}^{11}\mbox{-H}), & \sigma(\mbox{C}^{11}\mbox{-H}), \\ \mbox{H}) \rightarrow \sigma^*(\mbox{ring}), & \mbox{and } \sigma(\mbox{ring}) \rightarrow \sigma^*(\mbox{C}^{11}\mbox{-C}^{15}) & \mbox{charge} \end{array}$ transfer processes. On the other hand, the V_2' sine term reflects the steric and electrostatic destabilizing interactions between the positively charged methyl H_{21} atom and the aromatic H^7 or H^{10} , displaying energy maxima for C¹C⁶C¹¹C¹⁴ equal to -135° and 45° and energy minima for -45° and 135° .



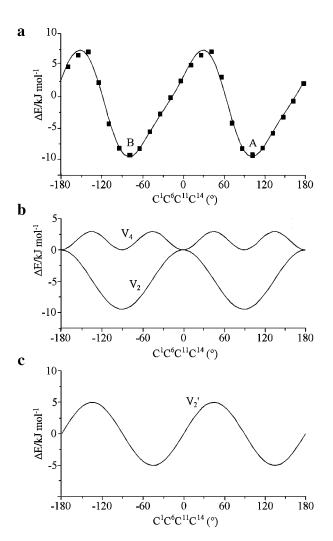


Figure 5. Optimized (B3LYP/6-31G^{*}) conformational energy profile for the internal rotation around the C⁶-C¹¹ bond of ibuprofen (a) and its Fourier deconvolution (b and c). $V_2 = -9.39$ kJmol⁻¹ and $V_4 = 2.87$ kJmol⁻¹ and $V'_2 = 5.27$ kJmol⁻¹.

The V_4 (2.87 kJmol⁻¹) Fourier term, in turn, represents the donor and acceptor hyperconjugative contributions, with respect to the π aromatic system, for each bond of the –CH₂C rotor: π (ring) \rightarrow $\sigma^{*}(C^{11}\text{-H}), \ \sigma(C^{11}\text{-H}) \rightarrow \pi^{*}(ring), \ \sigma(C^{11}\text{-C}^{15}) \rightarrow \pi^{*}$ (ring). These were found to contribute additively to the hyperconjugative effect of the whole group. In fact, while for $C^1C^6C^{11}C^{14}$ equal to $\pm 90^\circ$ the $\sigma(C^{11}\!-\!C^{15})\!\rightarrow\pi^*(ring)$ interplay is the predominant one, for the 0° and 180° conformations this effect occurs mainly through $\sigma(C^{11}-H) \rightarrow \pi^*(ring)$ interactions, the overall hyperconjugative stabilization being identical for the favored 0° , $\pm 90^{\circ}$, and 180° arrangements. From the magnitude of this term, as compared to V_2 and V'_2 , it is possible to conclude that hyperconjugation plays a less impor-

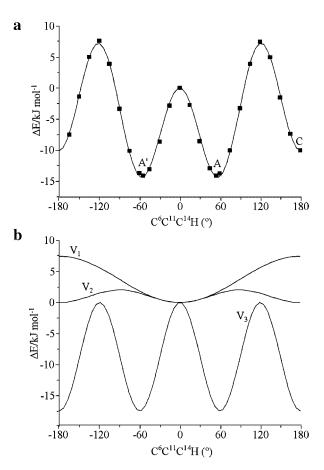


Figure 6. Optimized (B3LYP/6-31G^{*}) conformational energy profile for the internal rotation around the $C^{11}-C^{14}$ bond of ibuprofen (a) and its Fourier deconvolution (b). $V_1 = 7.53$ kJmol⁻¹, $V_2 = 2.07$ kJmol⁻¹, and $V_3 = -17.63$ kJmol⁻¹.

tant role than steric and electrostatic interactions, for the rotation of the isobutyl fragment within the ibuprofen molecule.

Rotation around the $C^{11}-C^{14}$

 C^{11} - C^{14} rotation interconverts conformers A, C, and A' (Fig. 6a), $C^6C^{11}C^{14}$ H being equal to 54.9°, 180.2°, and -54.4° , respectively. This rotation is practically symmetric around 0°. Actually, the presence of the *para*-substituent is responsible for a slight asymmetry which renders A and A' geometries not perfectly equivalent, although with identical conformational energies. The calculated value for the energy difference between conformers C and A, ΔE_{C-A} , is 4.12 kJmol⁻¹, the A \rightarrow C and A \rightarrow A' barriers being equal to 21.5 and 14.3 kJmol⁻¹, respectively (Fig. 6a).

Considering a Fourier decomposition of this energy variation profile, it was verified that only the cosine terms need to be considered (Fig. 6b). The highest contribution was found to come from V_3 (-17.63 kJmol⁻¹), thus favoring the staggered conformations, as expected for saturated hydrocarbons. The term in V_1 (7.53 kJmol⁻¹) favors the *cis* arrangement relative to the *trans*, since the concurrent steric repulsions between the two-methyl groups from the isobutyl moiety (particularly H¹⁷ and H²¹) and the phenyl π orbitals are expected to be very strong for this species. Thus, this contribution is responsible for the lower value of the A \rightarrow A' rotational barrier as compared to A \rightarrow C (Fig. 6a).

Vibrational Analysis

The ibuprofen Raman spectrum, in the 100–1750 and 2700–3300 cm⁻¹ regions, and the FTIR spectrum, in the 400–1750 and 2700–3300 cm⁻¹ intervals, for the solid state at room temperature, are represented in Figs. 7a and 7b, respectively. Table 3 comprises the experimental Raman and

FTIR wave numbers, as well as the *ab initio* calculated frequencies for the two most stable conformers, A and B (Fig. 1). A quite good accordance was found between the experimental and calculated values, after scaling according to Scott and Radom⁴⁰ in order to correct for the anharmonicity of the normal modes of vibration. Actually, since the calculated energy difference between B and the most stable A species is as small as 0.05 kJ mol^{-1} , both these conformers may have significative populations at room temperature. Nevertheless, the calculated values for the most conformationally sensitive frequency region (below 600 cm^{-1}) are remarkably consistent with the presence of conformer A alone. The bands at about 477, 523, and 586 cm^{-1} , for instance, display an outstandingly good agreement with the calculated values for A conformer (472, 512, and 597 cm^{-1}) when compared to the ones calculated for the B geometry.

Table 3 also contains a complete assignment of the ibuprofen observed FTIR and Raman bands to

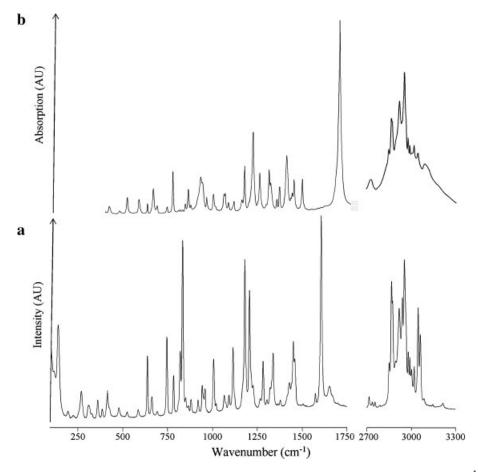


Figure 7. Vibrational optical spectra for solid ibuprofen. (a) Raman (100–1750 cm⁻¹, 2700–3300 cm⁻¹); (b) FTIR (400–1750 cm⁻¹, 2700–3300 cm⁻¹, in a KBr disk).

		Calcu	$lated^a$	
Raman	FTIR	Conformer A	Conformer B	Approximate $descriptions^b$
138		106 (3; 0)	117 (3; 0)	$C^{6}-C^{11}-C^{14}$, $C^{3}-C^{24}-C^{30}$ deformations
		163 (0; 0)	148 (0; 0)	$C^3-C^{24}-C^{26}$ deformation
193		216 (0; 1)	208 (0; 1)	$C^{26}H_3$ -C, $C^{19}H_3$ -C torsions
222		226 (1; 0)	224 (0; 0)	$C^{15}H_3-C$, $C^{19}H_3-C$ torsions
		231 (0; 1)	235(0; 1)	$C^{26}H_3$ –C torsion
		234(1; 0)	244(2;0)	$C^{19}H_3$ –C torsion; C^3 – C^{24} – C^{26} deformation
267		255 (0; 0)	250 (0; 0)	$C^{15}H_3$ –C, $C^{19}H_3$ –C torsions; C^3 – C^{24} – C^{26} deformation
306		299(1; 1)	288(1; 2)	$C^{11}-C^{14}-C^{15}$, $C^{26}-C^{24}-C^{3}$ deformations
312		200 (1, 1)	200 (1, 2)	
327		306 (1; 1)	319 (1; 0)	C^{11} – C^{14} – C^{15} deformation
359		353(2; 1)	350(2; 1)	C^{11} - C^{14} - C^{19} deformation; C^{11} - ϕ - C^{24} wagging (10b)
384		379(1; 0)	380 (0; 0)	$C^{11}-\phi-C^{24}$ in-plane out-of-phase bending (9b)
413		405 (3; 3)	409 (3; 1)	$C^{15}-C^{14}-C^{19}$ deformation
$413 \\ 424$	423	405 (0; 0)	409(3, 1) 424(1; 1)	ϕ out-of-plane bend (16a); C ¹⁴ (C ¹¹ C ¹⁵ C ¹⁹) sym. deform.
424	425	420 (0; 0)		$C^{24}-C^{30}$ -OH deformation; $C^{14}(C^{11}C^{15}C^{19})$ sym. deform.
477	400	479 (1, 10)	445 (1; 7)	$C^{24}-C^{30}$ -OH deformation; C (C C C) sym. deform. $C^{24}-C^{30}$ -OH deformation; ϕ in-plane bending (6b);
477	480	472 (1; 10)		
523	522	512 (1; 2)		CO–H bending; ϕ out-of-plane bending (16b)
			554 (1; 5)	CO-H bending; ϕ out-of-plane bending (16b)
FOR	500	FOF (0. 0F)	566 (1; 28)	ϕ out-of-plane bending (16b); C ²⁴ –C ³⁰ –OH deformation
586	589	597 (2; 35)		ϕ out-of-plane bending (16b); C ²⁴ –C ³⁰ –OH deformation
		605 (3; 47)	606 (3; 38)	CO-H out-of-plane bending; $C^{24}C^{30}=0$ def.; $C^{26}C^{24}C^{30}$ def.
637	636			
		623 (6; 30)	626 (6; 44)	CO–H out-of-plane bending; ϕ in-plane bending (6b)
662	669			
		638 (4; 13)	640 (4; 23)	CO-H out-of-plane bend; $C^{24}C^{30}=0$ def.; ϕ in-pl. bend (6b)
692	691	690 (1; 51)	694 (1; 39)	C–OH stretch; C^3 – C^{24} stretch; ϕ out-of-plane bend (4)
		726 (9; 8)	719 (4; 16)	ϕ out-of-plane bending (4); C=O out-of-plane wagging
746	747			
		767 (5; 16)	776 (21; 6)	CH_3 rocking; ϕ CH out-of-plane bending (17b)
783	780			
		792 (8; 11)	787 (1; 14)	φ out-of-plane bend (4); CH ₃ rock; C=O out-of-plane wagg.
809	810			
820	820			
		808 (16; 1)	805 (6; 1)	$C^{26}H_3$ rocking; C^6-C^{11} stretch.; $C^{14}(C^{11}C^{15}C^{19})$ sym. stretch
833	834			
		820 (4; 4)	824 (7; 2)	φ CH out-of-plane bending (10a)
850	850			
		827 (5; 2)	827 (7; 3)	ϕ CH out-of-plane bending (10a); C ²⁶ H ₃ rocking
865	867	839 (2; 14)	836 (4; 17)	ϕ CH out-of-plane bending (17b); CH ₃ rocking
880	880	866 (4; 2)	866 (3; 2)	CH_2 rocking; $C^{19}H_3$ rocking; $C^{14}H^{20}$ bending
		922 (2; 1)	921 (4; 1)	H ⁹ –C ⁴ C ⁵ –H ¹⁰ out-of-phase bending; C ¹⁵ H ₃ rocking
920	${\sim}921$			
		927 (7; 0)	928 (5; 0)	ϕ CH out-of-plane bend. (17a); $C^{15}H_3$ rocking
	937			(CO-H bending (H-bonded))
943	945	937 (5; 0)	938 (3; 0)	$H^7-C^1C^2-H^8$ out-of-phase bend; $C^{15}H_3$ rock; $C^{19}H_3$ rock
-		~~, ~/	×-, -,	(CO–H bending (H-bonded))
951		941 (2; 0)	939 (3; 0)	H^7 – C^1C^2 – H^8 out-of-phase bend; $C^{15}H_3$ rock; $C^{19}H_3$ rock
959	970	975 (9; 3)	975 (8; 3)	$C^{26}H_3$ rocking; $C^{26}-C^{24}-C^{30}$ antisymmetric stretching
1007	1008	1002 (0; 5)	1001 (0; 5)	ϕ CH in-plane bend. (18a)
1021	1020	(0,0)		· · · · · · · · · · · · · · · · · · ·
		1050 (2; 39)	1051 (2; 33)	CH ₂ twisting; C ²⁶ H ₃ rocking; C ²⁴ –H ²⁵ bend
		1000 (1,00)	1001 (1,00)	
				(Continued

 $\label{eq:constraint} \textbf{Table 3.} \quad \text{Experimental (Raman and FTIR) and DFT-Calculated Harmonic wave numbers (cm^{-1}) and Intensities for Ibuprofen$

(Continued)

Table	3.	(Continued)
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		Calcu	$lated^a$	
Raman	FTIR	Conformer A	Conformer B	Approximate descriptions ^{b}
1066	1068	1061 (1; 16)	1062 (1; 7)	CH_2 twisting; $C^{26}H_3$ rocking; C^{24} – H^{25} bend
1074	1074	1072 (3; 19)	1072 (3; 32)	CH_2 twisting; $C^{26}H_3$ rocking; $C^{19}H_3$ rocking; $C^{14}-H^{20}$ bend
1093	1092	1101 (12; 1)	1101 (9; 1)	$C^{15}H_3$ rocking; $C^{14}(C^{11}C^{15}C^{19})$ antisymmetric stretching
1115	~ 1116	1113 (1; 21)	1113 (4; 13)	φ CH in-plane bending (18b)
~ 1124	1123	1134 (1; 215)	1134 (1; 208)	C–O stretching; (CO–H in-plane bend)
		1157 (4; 10)	1157 (4; 12)	$C^{15}H_3$ rocking; $C^{19}H_3$ rocking; $C^{11}-C^{14}$ stretching
${\sim}1168$	1168	1171(42; 16)	1171 (37; 13)	ϕ CH in-plane bending (9a); C ³ –C ²⁴ stretching
~1100	1100	1179 (9; 10)	1180 (12; 9)	ϕ CH in-plane bending (9a); C ³ -C ²⁴ stretching
1181	1184	1170 (0, 10)	1100 (12, 5)	ϕ off in-plane behaving (5a), $\phi = \phi$ - stretching
1101	1104	1189 (13; 2)	1188 (14; 3)	C_6 – C_{11} stretching
1207	1208	1211 (6; 1)	1211(6; 2)	$C_6 - C_{11}$ stretching CH_2 twisting
1207 1227	$\frac{1208}{1231}$	1211(0, 1)	1211(0, 2)	CO–H in-plane bending (H-bonded)
1227	1201	1946 (C. E)	1945 (C. E)	C^{24} -H ²⁵ bend; (CO-H in-plane bending)
1007	1960	$1246 \ (6; 5)$	1245 (6; 5) 1272 (0; 2)	$C = \Pi$ bend; (CO = \Pi In-plane bending) $C^{24} = \Pi^{25}$ has diver. (CO = II in plane has diver.)
1267	1269	1274(7; 3)	1273 (9; 3)	C^{24} -H ²⁵ bending; (CO-H in-plane bending)
1283	${\sim}1283$	1278 (20; 4)	1281 (16; 1)	CH_2 wagging; C^{24} – H^{25} bending
1305	1306	1308 (3; 3)	1308 (2; 2)	ϕ in-plane bend (3); CH ₂ twisting; C ¹⁴ –H ²⁰ bending
1323	1321	1322 (2; 0)	1322 (5; 0)	ϕ in-plane bend (3); C^{24} – H^{25} bend + CO–H in-plane bend
${\sim}1332$	1330	1332 (3; 1)	1333 (2; 1)	C^{14} - H^{20} bending; C^{11} - H^{12} bending
1340		1336 (32; 4)	1339 (32; 4)	$C^{14}-H^{20}$ bending; $C^{11}-H^{13}$ bending
1366	1365	1361 (3; 55)	1361 (3; 52)	C^{24} -H ²⁵ bending; C^{24} -C ³⁰ -O ³² stretching; (CO-H bending)
		1375 (6; 4)	1374 (6; 4)	${ m C}_{ m sc}^{15}{ m H}_3$ and ${ m C}^{19}{ m H}_3$ out-of-phase symmetric deformation
1380	1380	1382 (3; 7)	1382 (3; 6)	$C_{15}^{26}H_3$ symmetric deformation
		1393 (2; 4)	1392 (2; 5)	$\mathrm{C^{15}H_3}$ and $\mathrm{C^{19}H_3}$ in-phase symmetric deformation
${\sim}1418$	1421	1413 (0; 10)	1413 (6; 1)	φ C–C stretching (19b); CH $_2$ twist; C ³ –C ²⁴ –H ²⁵ deformation
1430	1421			CO–H bending (H-bonded)
1443	1443	1453 (7; 1)	1453 (34; 1)	CH_2 deformation; $C^{15}H_3$ and $C^{19}H_3$ antisymmetric deformation
1452	1452	1460 (35; 1)	1461 (16; 4)	${ m C}^{15}{ m H}_3$ and ${ m C}^{19}{ m H}_3$ antisymmetric deformation
		1462 (14; 4)	1463 (2; 3)	$C^{26}H_3$ antisymmetric deformation
		1467 (1; 2)	1466 (2;3)	CH_2 deformation; $C^{15}H_3$ and $C^{19}H_3$ antisymmetric
				deformation
1460	1462	1469 (16; 8)	1469 (18; 7)	$C^{26}H_3$ antisymmetric deformation
1100	110-	1479 (33; 3)	1475(34; 3)	$C^{15}H_3$ and $C^{19}H_3$ in-phase antisymmetric deformation
		1481 (1; 9)	1481 (1; 9)	$C^{15}H_3$ and $C^{19}H_3$ in-phase antisymmetric deformation
1508	1508	1501 (1; 22)	1501 (0; 22)	ϕ C–C stretching (19a); C ³ –C ²⁴ and C ⁶ –C ¹¹ stretching
1575	1000	1566 (4; 0)	1565 (3; 0)	ϕ C–C stretching (9b)
1608		1606 (100; 1)	1606 (100; 1)	ϕ C–C stretching (9a)
1654		-	-	Overtone $(2 \times 833 \text{ cm}^{-1})$
1004	1720	1765 (6; 230)	1764 (6; 230)	$C^{30}=O^{31}$ stretching (H-bonded)
2717	1720	1705 (0, 250)	1704(0, 200)	Combination mode
2717 2737	2731	-	-	Combination mode
	2131	_	_	
2755		_	_	Combination mode
2787	0051	-	-	Combination mode
2851	2851			C^{14} II ²⁰ standaling
2867	2869	2899 (92; 7)	2900 (92; 6)	C^{14} -H ²⁰ stretching
2874	2874	2911 (35; 27)	2911 (36; 29)	$ m CH_2$ symmetric stretching
2895		0010 (01 05)		
	a	2918 (21; 36)	2918 (9; 30)	$\mathrm{C^{15}H_3}$ and $\mathrm{C^{19}H_3}$ out-of-phase symmetric stretchings
2919	2923	2924 (231; 37)	2924 (240; 42)	$C^{15}H_3$ and $C^{19}H_3$ in-phase symmetric stretchings

(Continued)

		Calcu	$lated^a$	
Raman	FTIR	Conformer A	Conformer B	Approximate descriptions ^{b}
2941		2945 (141; 26)	2944 (145; 28)	C ²⁶ H ₃ symmetric stretching
		2949 (28; 20)	2949 (28; 21)	CH ₂ antisymmetric stretching
2955	2956			
		2962 (57; 10)	2963 (56; 10)	C^{24} – H^{25} stretching
${\sim}2965$	${\sim}2965$			
		2978 (17; 8)	2978 (17; 8)	$\mathrm{C^{15}H_3}$ and $\mathrm{C^{19}H_3}$ out-of-phase antisymmetric stretchings
2981	2980			15 10
		2984 (90; 58)	2984 (97; 62)	C ¹⁵ H ₃ and C ¹⁹ H ₃ in-phase antisymmetric stretchings
2992	2992	2989 (116; 48)	2990 (106; 44)	C ¹⁵ H ₃ antisymmetric stretching
3003		3003 (55; 37)	3003 (60; 36)	C ¹⁹ H ₃ antisymmetric stretching
		3010 (104; 30)	3010 (109; 29)	C ²⁶ H ₃ antisymmetric stretching
3021	3020			22
		3031 (31; 14)	3030 (32; 14)	C ²⁶ H ₃ antisymmetric stretching
		3048 (50; 14)	3050 (38; 16)	H ⁹ –C ⁴ C ⁵ –H ¹⁰ antisymmetric stretching
3047	3045			
		3055 (50; 16)	3052 (67; 13)	H ⁷ –C ¹ C ² –H ⁸ antisymmetric stretching
		3065 (119; 21)	3067 (112; 23)	H ⁹ –C ⁴ C ⁵ –H ¹⁰ symmetric stretching
3063				
		3079 (92; 5)	3078 (94; 5)	H ⁷ –C ¹ C ² –H ⁸ symmetric stretching
3087				
3146		-	-	Overtone/combination mode
3215		-	-	Overtone/combination mode
		3536 (169; 47)	3537 (156; 46)	O ³² –H ³³ stretching

^{*a*}At the B3LYP/6-31G^{*} level of calculation; wave numbers above 600 cm⁻¹ scaled by a factor of 0.9614 ((Scott and Radom (40)); in parentheses: Raman scattering activities in $\mathbb{A} \cdot \mathbf{u}^{-1}$ and infrared intensities in km mol⁻¹.

^bThe commonly used Wilson notation for descriptions of benzene derivatives normal vibrations (Wilson Jr.⁴¹); Varsányi 1974⁴²) is presented inside parentheses.

the corresponding normal modes of vibration. This assignment was based both in the calculated wave numbers and intensities, despite the latter are recognizably much less accurate. It may be concluded that the experimental vibrational spectra confidently reflect the presence of specific intermolecular interactions involving the carboxylic group. In fact, the band at 1720 cm^{-1} , assigned to the $C^{30} = O^{31}$ stretching, displays a downward shift relative to the calculated value for the isolated molecule (Tab. 3). A similar deviation to low wave numbers is suggested to occur for the O-H stretching vibration, since this mode is not observed at the predicted 3536 cm^{-1} frequency, probably occurring as a broad band at about 2900-3000 cm⁻¹, overruled by the intense C–H stretching features. The CO-H bending modes, in turn, display clear upward shifts relative to the calculated values for isolated ibuprofen: the calculated modes at 1123, 1246, and 1361 cm^{-1} corresponding to the experimental 1231, 1322, and 1430 cm^{-1} frequencies, respectively (Tab. 3). Moreover, the 937 and 945 cm^{-1} features, for

instance, display a particularly high intensity in FTIR (as opposed to Raman), thus being empirically ascribed to CO–H bending modes. This spectral behavior is characteristic of hydrogen bond type interactions. These close-contacts, which occur through the O^{31} and H^{33} atoms within the ibuprofen molecule, are also responsible for the presence of ibuprofen dimeric entities in condensed phases.³⁵

CONCLUSIONS

The present study reports a thorough conformational analysis of the ibuprofen molecule, by DFT based methods. Eight different energy minima were found, displaying an s-*cis* orientation of the carboxylic group. The s-*trans* arrangement was not considered, since it was previously verified, for the analogous molecules, to be significantly less stable than its s-*cis* counterpart.

It was found that the relative orientation of the substituents in the ibuprofen molecule, which can

be considered as a *para*-substituted phenyl ring, hardly affects its conformational stability. These groups can be located below or above the ring plane, either to the same side or in opposite positions, leading to only four low energy sets of conformers: AB ($\Delta E \approx 0$ kJmol⁻¹), CD ($\Delta E \approx 4.1$ kJmol⁻¹), EF ($\Delta E \approx 5.2$ kJmol⁻¹), and GH ($\Delta E \approx 9.2$ kJmol⁻¹), with populations at room temperature of 75, 14, 9, and 2%, respectively (Tab. 1). When considering Gibbs populations, however, these change to 84% ($\Delta G \approx 0$ kJmol⁻¹), 5% ($\Delta G \approx 6.7$ kJmol⁻¹), 10% ($\Delta G \approx 5.4$ kJmol⁻¹), and 1% ($\Delta G \approx 11.8$ kJmol⁻¹), due to the inclusion of the entropy effect.

Despite the very small energy difference between the A and B conformations, the optical vibrational spectroscopy data reflects the sole presence of A in the solid phase, most probably on account of a better packing of this structure in the crystal. Furthermore, spectroscopic evidence of intermolecular hydrogen bonds between the carboxylic groups of adjacent ibuprofen molecules, leading to the formation of dimers (in the condensed phase) was also obtained. Disruption of these dimeric species may be one of the key factors for the understanding of the pharmacokinetics and drug release behavior of ibuprofen containing controlled release systems.

This kind of conformational analysis based on both spectroscopic and theoretical studies yields information of the utmost relevance for the elucidation of the SAR's ruling the pharmacological characteristics of substituted α -arylpropionic acids, mainly when coupled to biochemical activity evaluation experiments.

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