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Short induction gap gas electron multiplier (GEM) for X-ray spectroscopy

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Abstract

Experimental work was carried out to evaluate the performance of a Gas Electron Multiplier (GEM) operated with a micromesh readout plane that enabled the induction gap to be set at $50\,\mu m$. We measured the essential operational parameters of this system using Ar(75%)-isobutane (25%) as the counter gas mixture. The measurements included the effective gain, effective gain stability, and the X-ray energy resolution using a $5.89\,keV$ X-ray source. These studies demonstrated several advantages of the current system when compared with the standard operation, such as lower operational voltages, higher effective gains and improved effective gain stability. \odot 2006 Elsevier B.V. All rights reserved.

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1. Introduction

A typical Gas Electron Multiplier (GEM) configuration [1] sets the induction gap at 1 mm or more where the induction gap is defined as the distance between the bottom GEM electrode and the readout plane. It is well known that Kapton-based GEMs are susceptible to absorption by water vapour and other counter gases [2]. One of the consequences of the absorption process is the sagging of the GEM foil thereby changing the induction gap from its initial set value. However, the effective gain of a GEM is strongly dependant upon the induction field [3,4]. As a result, GEM sagging leads to an effective gain instability [4]. One way to circumvent this is to introduce dielectric pillars between the GEM foil and the readout plane at regular interval [5]. In the present study, we have used a standard GEM that was coupled with a micromesh [6] readout plane with 50 µm tall Kapton pillars patterned on the micromesh at 2 mm pitch. The 50 µm induction gap had several distinct operational advantages. For example, the effective gain stability was improved owing to a good induction gap definition and the absolute voltage needed to sustain a particular induction field was lowered by a factor of 20.

2. Method

The X-ray sensitive area of the present detection system consisted of a $10\,\mathrm{mm} \times 10\,\mathrm{mm}$ GEM, a micromesh and a $100\,\mu\mathrm{m}$ thick aluminium foil that defined the drift window. The drift distance was set at 5 mm whereas the induction gap was set at 50 $\mu\mathrm{m}$. The GEM used here was of a standard geometry and was fabricated at the CERN TS-DEM workshop and consisted of a 50 $\mu\mathrm{m}$ thick copper clad (5 $\mu\mathrm{m}$) Kapton foil with 70 (50) $\mu\mathrm{m}$ holes patterned at 140 $\mu\mathrm{m}$ hole pitch. The micromesh was also manufactured at CERN TS-DEM workshop and consisted of a 5 $\mu\mathrm{m}$ thick Copper mesh with 25 $\mu\mathrm{m}$ holes etched at a pitch of 50 $\mu\mathrm{m}$ with 50 $\mu\mathrm{m}$ tall and 150 $\mu\mathrm{m}$ diameter Kapton pillars distributed at 2 $\mu\mathrm{m}$ mintervals [6].

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In the experimental studies described in the following sections, a Mn-K X-ray (5.89 keV) beam illuminated the detector drift space perpendicular to the GEM and micromesh planes through an entrance window near the end wall of the detector chamber. The entrance window consisted of 100 µm mylar mesh with 100 nm aluminium coating on one side. The entrance window was located approximately 10 mm away from the drift window. In this study, the detection chamber was operated at a constant gas flow rate where the gas composition was controlled by Brooks mass-flow controllers (model 5850E) built into a rig that was constructed using stainless-steel tubing. An argon (75%) and isobutane (25%) counter gas mixture was used throughout this study. The ambient pressure and the temperature of the detector box were recorded with every measurement. The global count rates used throughout these studies were kept in the region of 10 kHz.

The drift electrode and the GEM mesh were operated negative with respect to the micromesh that was held close to earth. The micromesh was connected electrically to an Ortec preamplifier (model 142A). The preamplifier output was then fed into an Oretc shaping amplifier (model 575A) with shaping time constants adjusted to 0.5 µs. The bipolar output of the shaping amplifier was in turn fed into an Ortec pulse height analyser (Ortec Trump-PCI-2 K plug in card with Maestro-32 software for Windows). The effective gain and the X-ray energy resolution were examined as a function of the induction field, $E_{\rm I}$, and the voltage differences applied across the GEM holes, $\Delta V_{\rm GEM}$. The drift field, E_d , was maintained at approximately 3.5 kV/cm throughout these studies. The effective gain stability was also studied by monitoring the detector gain for a period of one month.

3. Effective gain and X-ray energy resolution

Fig. 1 shows the variation of the effective gain as a function of the induction field for the range $E_{\rm I}=0.4$ –40 kV/cm for a number of different voltages across the GEM holes ($\Delta V_{\rm GEM}=400$, 450 and 500 V). The effective gain increases almost linearly with increasing induction field. When $E_{\rm I}$ exceeds values above 15 kV/cm, further electron multiplication begins in the induction region (parallel plate amplification mode).

Fig. 2 shows the variation of the X-ray energy resolution at $5.89 \, \mathrm{keV}$ as a function of the induction field for different voltages across the GEM holes. For induction fields, E_{I} , lower than $2 \, \mathrm{kV/cm}$, the X-ray energy resolution was rather poor due to most electrons being directed to the lower GEM electrode. The optimum X-ray energy resolution was observed in the induction field region $2-15 \, \mathrm{kV/cm}$, beyond which rapid deterioration in the X-ray energy resolutions occurred.

Fig. 3 shows the variation of the effective gain as a function of voltage applied across the GEM holes for a number of different induction fields ($E_I = 6$, 12, 20, 30 and

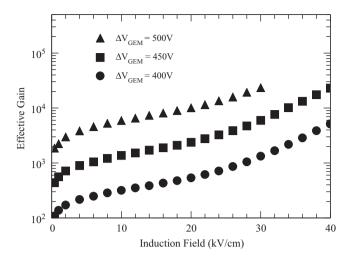


Fig. 1. A plot of the effective gain as a function of the induction field using Ar(75%)-isobutane(25%) for $\Delta V_{\rm GEM}=400,\,450$ and 500 V. In all cases the drift field $E_{\rm d}=-3.5\,{\rm kV/cm}.$

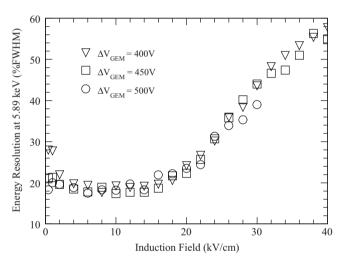


Fig. 2. X-ray energy resolution (% FWHM) of the Mn K X-rays (5.89 keV) as a function of the induction field using Ar (75%)-isobutane (25%) for $\Delta V_{\rm GEM} = 400$, 450 and 500 V. In all cases $E_{\rm d} = -3.5\,{\rm kV/cm}$.

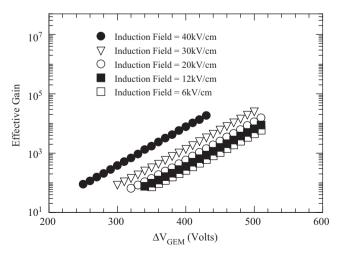


Fig. 3. Effective gain as a function of the potential across the GEM holes ($\Delta V_{\rm GEM}$) using Ar (75%)-isobutane (25%) for a number of different induction fields. In all cases the drift field $E_{\rm d}=-3.5\,{\rm kV/cm}$.

 $40\,\mathrm{kV/cm}$). The highest effective gain at the lowest induction field used here was found to be approximately 6000. Increasing the induction fields to values higher than $6\,\mathrm{kV/cm}$ shifted the effective gain-voltage (ΔV_{GEM}) characteristics. The highest gain observed was approximately 27,000 when the induction field was set at $30\,\mathrm{kV/cm}$.

The X-ray energy resolution at 5.89 keV associated with each curve in Fig. 3 is plotted in Fig. 4. It is clear from Figs. 1–4 that although effective gain increased with induction fields higher than $6 \, \text{kV/cm}$, this was accompanied by a degradation in the corresponding X-ray energy resolution [5]. Fig. 5 shows an optimum pulse height spectrum at 5.89 keV when the GEM was operated at an effective gain of 1900 ($\Delta V_{\text{GEM}} = 470 \, \text{V}$, $E_{\text{I}} = 6 \, \text{kV/cm}$, $E_{\text{d}} = -3.5 \, \text{kV/cm}$). The X-ray energy resolution derived from this figure was found to be 15.2% FWHM.

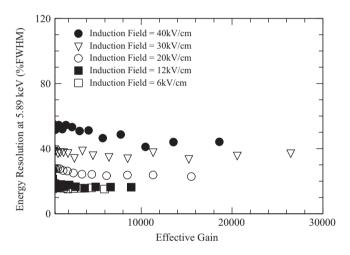


Fig. 4. X-ray energy resolution (%FWHM) of the Mn K X-rays (5.89 keV) as a function of the effective gain using Ar (75%)-isobutane (25%) for a number of different induction fields. In all cases the drift field $E_{\rm d}=-3.5\,{\rm kV/cm}$.

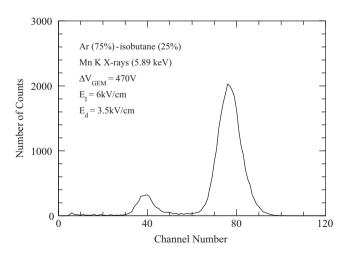


Fig. 5. A typical pulse height spectrum for 5.89 keV X-rays using the GEM at an effective gain of 1900 in Ar (75%)-isobutane (25%) with $\Delta V_{\rm GEM}=470\,{\rm V},\,E_{\rm I}=6\,{\rm kV/cm}$ and $E_{\rm d}=-3.5\,{\rm kV/cm}.$

4. Effective gain stability

The effective gain of flow type gaseous detectors is influenced by the ratio P/T [7 and references therein], where P is the gas pressure in mB and T is its absolute temperature in K. Fig. 6 shows the effective gain of the present GEM over a one month period plotted as a function of P/T. Fitting an exponential curve to the data gives the effective gain sensitivity of the detector to changes in the ambient conditions. In the present case, this evaluates to $1.55 \, \text{K/mB}$ as shown in Fig. 6. Assuming that P is roughly constant at around $1000 \, \text{mB}$ and allowing for maximum temperature excursion of $\pm 5 \,^{\circ}\text{C}$ from $20 \,^{\circ}\text{C}$ would result in a maximum gain excursion of $\pm 8.9\%$.

5. Conclusion

A short gap GEM was successfully operated with an induction gap of $50\,\mu m$ to investigate parameters relevant to the X-ray detection. These included measurements of the effective gain and X-ray energy resolution as a function of the induction field and the voltage applied across the GEM holes.

The highest effective gain was found to be approximately 27 000 when the voltage across the GEM holes was set at 500 V and induction field at 30 kV/cm. However, the best X-ray energy resolution was found to be 15.2% FWHM when the voltage across the GEM holes was 470 V (effective gain 1900) and the induction field was 6 kV/cm.

The effective gain shift with time experienced during our previous studies [4] was overcome by the simple introduction of $50\,\mu m$ tall Kapton pillars between the lower GEM electrode and the readout plane. Effective gain shifts observed during the one month run were consistent with the changes in ambient pressure and temperature over that period.

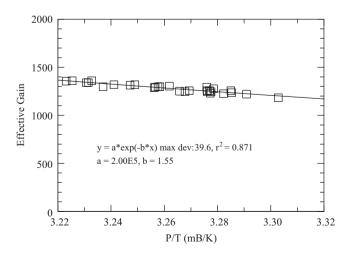


Fig. 6. The effective gain measured over a period of one month using Ar (75%)-isobutane (25%) as a function of the ambient parameter P/T. $\Delta V_{\rm GEM} = 450 \, {\rm V}$, $E_{\rm I} = 6 \, {\rm kV/cm}$ and $E_{\rm d} = -3.5 \, {\rm kV/cm}$.

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