A CFD based numerical study on aerodynamic characteristics of π cross sections using baffle plates

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Abstract

The aerodynamic instability of bridges should be one of the biggest concerns for bridge designers. From all the aerodynamic studies, only a few are related to improvements of bridge cross sections. Among them, strategies such as grating, edge plates, edge fairing plates, side plates, baffle plates or flaps have been tested. The aim of this study is associated with the efficiency of using baffle plates with the purpose of improving the aerodynamic characteristics of a π cross section (B/D=6). The limit will be the use of a rectangular cross section. The Scanlan model, namely the A_2^* coefficient, is considered with the aim of evaluating the aerodynamic efficiency of the cross section. In order to determine the fluid flow around the obstacle, a numerical algorithm of Computational Fluid Dynamics (CFD) based on the Finite Volume Method (FVM) is used. Additionally, the Forced Oscillation Method (FOM) is adopted for evaluating aeroelastic coefficients.

Keywords: CFD, FVM, FOM, bridges, Scanlan model, aerodynamic coefficients, flutter, improving cross section.

1 Introduction

Wind action is one of the most determining factors for the safety of large and flexible structures. As it is well known, since the famous Tacoma Narrows Bridge failure, in 1940, the design of long span cable-stayed and suspension bridges requires careful study of their aerodynamic behaviour under wind loads.



Traditionally, the characterisation of aerodynamic wind action and its effects on flexible structures have been based on physical models tested in wind tunnels. More recently, an alternative numerical approach has been developed and refined [1,2]. This empirical theory, based on the so-called Scanlan model for the evaluation of wind forces also called aeroelastic forces (AF), involves important simplifications. However, this numerical approach requires the identification of several coefficients whose estimation assumes then central importance in the evaluation of the response of long-span bridges to wind loading. These aerodynamic coefficients strongly depend on the bridge cross section and on a particular dimensionless velocity (reduced velocity).

On the other hand, cable stayed bridge girders with two I-beams have been adopted in long span bridges because their structural and economical advantages. However, this basic π cross section does not necessarily have good aerodynamic stability. In some cases, depending on bridges characteristics, some aerodynamic or structural measures against wind-induced vibrations are needed. Fairings, flaps, edge plates, side plates, baffle plates or gratings are usually used to aerodynamically improve this cross section. These attachments can be added to the structure during the construction or after that. Active or passive control of aerodynamic characteristics is not usually used as aerodynamic stability solution of a π cross section.

The intention of this study is associated with the efficiency of using baffle plates with the purpose of improving the aerodynamic characteristics of a π cross section. Accordingly, it is considered 5 sections: the π section without any baffle plate as a basic section, 3 π sections with 1, 2 or 3 baffle plates, and the rectangular section. It is consider the Scanlan model, namely the A_2^* coefficient, with the aim of evaluating the aerodynamic efficiency of the cross section.

In order to determine the fluid flow around the obstacle, it is used a numerical algorithm of CFD based on the FVM. The implemented program is suitable to simulate incompressible and isothermal bidimensional unsteady fluid flows around obstacles. It is assumed that the flow's domain may be discretised in a Cartesian and structured control volume mesh, whose faces have vertical and horizontal directions. In this algorithm, the high Reynolds number $k-\varepsilon$ turbulence diffusion model is applied to simulate the flow turbulence. All relevant equations used for modulation of fluid flow can be consulted in reference [3]. The obstacles movements were modelled indirectly by changing the velocity components of the fluid flow at external inlet boundary domain. Additionally, FOM is adopted for evaluating aeroelastic coefficients (AC).

2 Scanlan model

The term "flutter" was initially used by aeronautic engineers to describe the aerodynamic instability of aircraft wings, which is characterised by both vertical and torsional oscillations. Flexible structures, such as long-span bridges, under air fluid flow action also experiment similar unstable effects. In this case, the so-called "flutter" phenomenon happens when one or more oscillating modes show

increasing amplitudes due to aerodynamic forces, whose growth also depends on the structural movements. That is to say, aerodynamic forces, now called aeroelastic forces (AF), are dependent, not only on geometry of the cross section and on velocity of the free flow, but also on the structural movements and viceversa. In this domain, it is usual to call self-excited forces to AF.

The first analytical model to explicit those AF, in the field of Bridge Aerodynamics, was presented by Scanlan and Tomko [1], with the aim of analyzing the "flutter" stability. This model considered only two degrees of freedom (one vertical and one angular), where lift and moment AF were dependent on rotation and its velocity, and vertical velocity. On the other hand, this model was applied under the following assumptions: i-) the free flow had no oscillation; ii-) the movements had constant frequency; iii-) and the amplitude of movements was incipient. After that, this approach was improved to complete the model [4] which is presented below considering a particular case of a cross section that has three degrees of freedom, as indicated in fig. 1. In this model, it is assumed that any AF is dependent on all movements through the displacement and velocity components.

Accordingly, the dynamic system of balanced equations can be express by

$$\begin{cases} M_{1} \cdot \ddot{a}_{1} + C_{1} \cdot \dot{a}_{1} + K_{1} \cdot a_{1} = F_{a1} \\ M_{2} \cdot \ddot{a}_{2} + C_{2} \cdot \dot{a}_{2} + K_{2} \cdot a_{2} = F_{a2} \\ M_{12} \cdot \ddot{a}_{12} + C_{12} \cdot \dot{a}_{12} + K_{12} \cdot a_{12} = F_{a12} \end{cases}$$
 (1)

where M_i , C_i , K_i , a_i , \dot{a}_i and \ddot{a}_i correspond to the mass, damping, stiffness, displacement, velocity and acceleration of cross section according to direction i. Before progressing, it is worth mentioning two important dimensionless parameters: reduced velocity Ur and reduced frequency Kr, defined by

$$Ur = \frac{U}{fB} = \frac{2\pi}{Kr} \qquad ; \qquad Kr = \frac{B\omega}{U}$$
 (2)

where U is the velocity of the free fluid flow and $\omega = 2\pi f$ represents the angular frequency of the system oscillation. Dimension B is indicated in fig. 1.

According to the Scanlan model, the characterisation of $AF F_a$ is made by means of some constants, by AC and by movements of the cross section a, i.e., taking into consideration the definition of force coefficients it is possible to write down

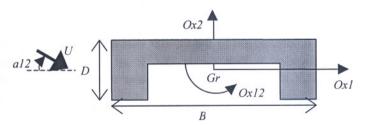


Figure 1: Degrees of freedom of a cross section.

$$F_{a} = \begin{bmatrix} F_{a1} \\ F_{a2} \\ F_{a12} \end{bmatrix} = CF \cdot PHA^{*} \cdot a \qquad \text{and} \quad CF = \frac{\rho U^{2} B}{2} \begin{bmatrix} 1 & 0 & 0 \\ 0 & 1 & 0 \\ 0 & 0 & B \end{bmatrix}$$
(3)

where CF is a matrix of force coefficient constants, ρ represents the density of the fluid, and

$$PHA^{*} = \begin{bmatrix} \frac{KrP_{5}^{*}}{U} & \frac{KrBP_{2}^{*}}{U} & Kr^{2}P_{3}^{*} & \frac{Kr^{2}P_{6}^{*}}{B} & \frac{KrP_{1}^{*}}{U} & \frac{Kr^{2}P_{4}^{*}}{B} \\ \frac{KrH_{1}^{*}}{U} & \frac{KrBH_{2}^{*}}{U} & Kr^{2}H_{3}^{*} & \frac{Kr^{2}H_{4}^{*}}{B} & \frac{KrH_{5}^{*}}{U} & \frac{Kr^{2}H_{6}^{*}}{B} \\ \frac{KrA_{1}^{*}}{U} & \frac{KrBA_{2}^{*}}{U} & Kr^{2}A_{3}^{*} & \frac{Kr^{2}A_{4}^{*}}{B} & \frac{KrA_{5}^{*}}{U} & \frac{Kr^{2}A_{6}^{*}}{B} \end{bmatrix}$$
(4)

$$a^{T} = \begin{bmatrix} \dot{a}_{2} & \dot{a}_{12} & a_{12} & a_{2} & \dot{a}_{1} & a_{1} \end{bmatrix}$$
 (5)

and the coefficients P_i^* , H_i^* e A_i^* are the AC, also called Scanlan coefficients or aerodynamic derivatives. Usually, these AC are presented as a function of only two factors: geometry of cross section and reduced velocity (or reduced frequency). It is said that this linear model is valid only for incipient amplitude of cross section movements, and also only when the frequency of oscillation is quite far from the Strouhal frequency. It is also known that the incoming flow and the amplitude of oscillation are very important on the evaluation of these coefficients in the experimental field, as well as in the methodology used to get them (forced oscillation or free oscillation methods).

Based on Fourier transform, the transformation of Scanlan eqn (3) to frequency domain can lead to

$$\widetilde{F}_{a1} = \frac{\rho U^2 B K r^2}{2} \left[\left(P_4^* + i P_1^* \right) \frac{\widetilde{a}_1}{B} + \left(P_3^* + i P_2^* \right) \widetilde{a}_{12} + \left(P_6^* + i P_5^* \right) \frac{\widetilde{a}_2}{B} \right] \tag{6}$$

$$\widetilde{F}_{a2} = \frac{\rho U^2 B K r^2}{2} \left[\left(H_4^* + i H_1^* \right) \frac{\widetilde{a}_2}{B} + \left(H_3^* + i H_2^* \right) \widetilde{a}_{12} + \left(H_6^* + i H_5^* \right) \frac{\widetilde{a}_1}{B} \right]$$
(7)

$$\widetilde{F}_{a12} = \frac{\rho U^2 B^2 K r^2}{2} \left[\left(A_4^* + i A_1^* \right) \frac{\widetilde{a}_2}{B} + \left(A_3^* + i A_2^* \right) \widetilde{a}_{12} + \left(A_6^* + i A_5^* \right) \frac{\widetilde{a}_1}{B} \right]$$
(8)

where \widetilde{F}_i and \widetilde{a}_i correspond to the Fourier transform of AF and movements according to direction i. That means that, if all parameters are known or specified, and if it is calculated both the Fourier transform of movements \widetilde{a}_i according to a particular direction i and the Fourier transforms of all AF \widetilde{F}_i , then the correspondent AC can be evaluated using the last three equations.

For instance, on FOM context, used by Nakamura [5], the movements according to a particular direction i are imposed, so it is only needed the calculation of the three AF and the subsequent Fourier transforms to evaluate the six matching Scanlan coefficients.

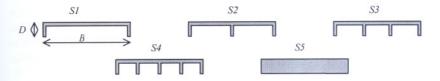


Figure 2: Five cross sections considered.

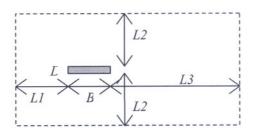


Figure 3: Main characteristics of the domain.

The displacement $a = a_0 \cos(\omega t)$ and the velocity $\dot{a} = -a_0 \omega \sin(\omega t)$ imposed to a cross section were defined by Nakamura as sinusoidal functions with constant amplitude a_0 and constant frequency ω . The symbol t represents the time.

3 Numerical procedures

To evaluate the efficiency of using baffle plates with the purpose of improving the aerodynamic characteristics of a π cross section (B/D=6), five different cross sections were used (see fig. 2): the π section without any baffle plate (SI) as a basic section, 3 π sections with 1 (S2), 2 (S3) or 3 (S4) baffle plates, and the rectangular section (S5). It is considered a structured control volume mesh whose main characteristics of the domain and the corresponding discretisation, according to fig. 3, are: B=12m; D=2m; L1=19.8m; L2=15.0m; L3=31.8m; maximum dimension=750mm; minimum dimension=51mm. The numbers of the Control Volumes used are from 81 to 89 for direction 2, and 167 to 227 for direction 1.

Usually, AC are presented graphically conditional on reduced velocity Ur. Accordingly, for each graph, it is used 9 points, from Ur=1.0 to Ur=1.5.

It is considered two velocities of the free fluid (air at standard conditions) flow: U=0.376m/s (Re=5E4) and U=37.6m/s (Re=5E6). The used time increment were $\Delta t=2E-2s$ and $\Delta t=2E-4s$ respectively. For each case, table 1 shows some static relevant values, where CFi represents the average aerodynamic force coefficient according to direction i, Δ means the average amplitude of variation, and St corresponds to the Strouhal number. The time interval is established in order to get stability during the simulation, as much as necessary.

Static results permit to draw some initial conclusions: firstly, the results obtained for the different cross sections are quite different and they are closely associated with the separation-and-reattaching flow and the vortex development between baffle plates, whose characteristics are not presented in this paper; secondly, as the velocity of free fluid flow grows, the amplitude oscillation of aerodynamic forces becomes lower and lower; thirdly, one can understand that the number of baffle plate's limit is the rectangular cross section.

Table 1: Some relevant values got when the cross section is at rest.

		CF1	∆CF1	CF2	△CF2	CF12	ACF12	St
SI	Re=5E4	1.181	0.332	-0.428	0.236	-0.239	0.174	0.120
	Re=5E6	1.217	0.332	-0.411	0.197	-0.292	0.168	0.124
S2	Re=5E0	1.133	0.272	-0.344	0.241	-0.148	0.201	0.120
	Re=5E6	1.122	0.078	-0.203	0.120	-0.192	0.101	0.122
S3	Re=5E4	1.106	0.102	-0.200	0.206	-0.027	0.179	0.116
	Re=5E6	1.068	0.007	-0.085	0.041	-0.044	0.030	0.118
S4	Re=5E4	1.130	0.049	-0.092	0.200	0.049	0.169	0.116
	Re=5E6	1.071	0.004	-0.024	0.021	0.027	0.014	0.118
S5	Re=5E4	1.134	0.035	-0.005	0.121	0.004	0.131	0.106
	Re=5E6	1.094	0.002	-0.001	0.041	0.000	0.025	0.116

Taking into consideration the amplitude of forced oscillation, two important aspects mentioned in the literature [2,5] should be referred: firstly, the Scanlan model is valid only for incipient amplitude oscillations of the cross section; secondly, this amplitude is usually dimensionless by means of D, with values inferior to 0.1. However, taking into consideration the physical aspect revealed by eqns (6-8), it is possible to conclude that, for forced rotations, the amplitude of displacement is the key parameter to state the changing conditions, but for forced translations, the important parameter is the velocity of forced movement, not the displacement. So, in this work, the amplitude of forced rotations is assumed as $a_0 = 3\%(rad) = 1.72^{\circ}$ while, for forced translational cases, with an angular frequency w, $a_0 * w \approx 3\% * U \Rightarrow a_0 = 0.005 B * Ur$. Last statement means

that the maximum velocity of forced oscillation is limited to 3% of the flow velocity.

Each simulation has four phases to be completed. In the first one, the velocity of the free flow is increased to reach the desired value keeping the cross section fixed. Next, the time step is adapted if necessary. The objective consists in having an appropriate discretisation (not less than 500 time steps) of the forced oscillation period. In the third phase, it is expected that the corresponding AF reach a regular pattern of oscillation. At last, it is made the record for evaluation of the correspondent AC. The record must have more than 20 000 time steps, or at least 5 forced oscillation periods.

4 Results

This work has produced a large amount of results, from which it is only possible to present the most important graphs which support the conclusions written down below.

In order to inspect the dependence of the baffle plate's number on the evaluated values of the AC, it is presented in figs. 4-7 the most important AC H_1^* and A_2^* used for calculation of vertical and torsional stability.

With the intention of confronting the current results, it is presented in figs. 8-9 the results obtained experimentally by Matsumoto [6] and by this methodology considering rectangular sections. The dimensions considered are B/D=6, in this study, and B/D=5 and B/D=8 for the Matsumoto cases.

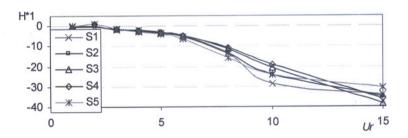


Figure 4: ACH_1^* evaluated for five different sections (Re=5E4).

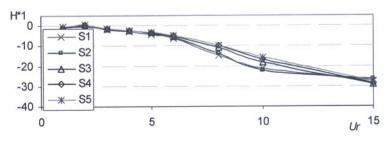


Figure 5: $AC H_1^*$ evaluated for five different sections (Re=5E6).



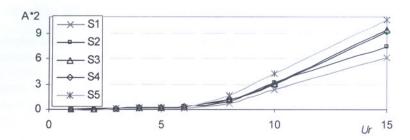


Figure 6: $AC A_2^*$ evaluated for five different sections (Re=5E4).

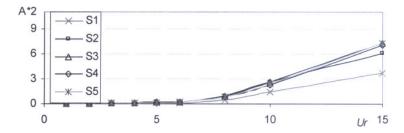


Figure 7: $AC A_2^*$ evaluated for five different sections (Re=5E6).

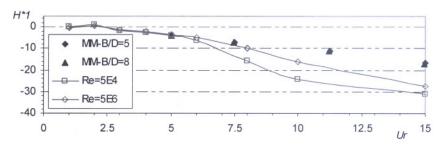


Figure 8: Comparing $AC H_1^*$ evaluated for rectangular sections.

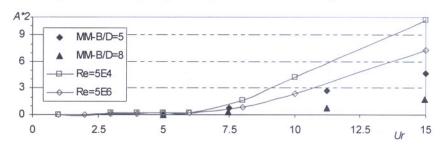


Figure 9: Comparing $AC A_2^*$ evaluated for rectangular sections.

The inspection of these results permit to draw the following particular conclusions, considering the $AC H_1^*$ and A_2^* :

- o the AC evaluated when considering moderate velocities (Re=5E4) diverge from those obtained when considering high convective flows (Re=5E6) (see figs. 4-7);
- o in particular, for the basic section SI, the AC A_2^* changes from 2.35 (Re=5E4) to 1.40 (Re=5E6) when reduced velocity stands for I0 (see figs. 6-7);
- o for example, considering one structure characterised by f=0.5Hz and with the stability limit of $A_2^*=2$, this could mean a critical flutter velocity of 28.8m/s (Ur=9.6 for Re=5E4 case) or 34.2m/s (Ur=11.4 for Re=5E6 case), which represents a significant difference;
- for AC H₁*, the differences evaluated for the five cross sections are not so important as the corresponding values are negative, i.e. they work for stability (see figs. 4-5);
- o for all the considered five sections, the torsional instability, i.e. the flutter phenomenon, will be less probable to happen for reduced velocities lower than 5, as the $AC A_2^*$ assumes quasi null values in this range;
- o on the other hand, when the reduced velocity is high, the most important AC A_2^* have significant values (see figs. 6-7);
- o taking into consideration only the changing of the $AC A_2^*$, one can bring up that the basic π section gets close to the rectangular section as the number of baffle plates increase;
- o considering only the possibility of torsional instability, i.e. the evaluation of $AC A_2^*$, figs. 6 and 7 show that the basic π section is also the less sensitive to flutter of the considered five sections;
- o taking into consideration the values obtained by Matsumoto (see figs. 8-9), it is possible to conclude that there is a poor agreement for Re=5E4 case, which is close to the velocity used by Matsumoto. Even, for Re=5E6 case, the corresponding results are not close, mainly for high reduced velocities range.

5 Conclusions

For evaluating the effects of wind action on flexible structures such as long-span bridges, it is generally used a numerical approach based on the so called Scanlan model, which requires the identification of some AC. The objective of this paper is to present a numerical study associated with the efficiency of using baffle plates with the purpose of improving the aerodynamic characteristics of a π cross section (B/D=6).

The results presented here are evaluated by a numerical approach based on an algorithm of *CFD* (*FVM*). The *FOM* is the methodology used for numerical evaluation of *AC*. The computer code developed on the basis of this



methodology is applied to the aeroelastic study of five cross-sections, which permits to characterise the influence of the baffle plate's number on the evaluation of the most important AC.

Considering the variation of the most important AC used for the analysis of vertical and torsional stability, one may essentially conclude that:

- o concerning vertical oscillations, the basic π section, including or not baffle plates, and the rectangular section are all stables;
- o taking into account the flutter phenomenon case, the results obtained suggest that the addition of baffle plates to a π section does not improve its aerodynamical characteristics;
- o the evaluation of the AC, namely A_2^* , depends on the velocity of the free flow considered.

Under these circumstances, it seems important to have specific rules in terms of the characteristics of incoming fluid flow in order to evaluate this important AC, whose values are determinant for the evaluation of the critical velocity of aeroelastic instability.

The conclusions drawn from the specific case of the π section can not be directly extrapolated to other examples, although the methodology presented in this paper can be applied to other cases with different shapes of cross section.

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