# RIEMANN-HILBERT PROBLEM ASSOCIATED WITH ANGELESCO SYSTEMS 

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#### Abstract

Angelesco systems of measures with Jacobi type weights are considered. For such systems, strong asymptotic development expressions for sequences of associated Hermite-Padé approximants are found. In the procedure, an approach from Riemann-Hilbert Problem plays a fundamental role.


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## 1. The statement of the Riemann-Hilbert problem

Let $\Delta_{j}=\left[c_{1, j}, c_{2, j}\right] \subset \mathbb{R}, j=1,2$, be two intervals which are symmetric with respect to the origin. This means that $c_{1,1}=-c_{2,2}$ and $c_{1,2}=-c_{2,1}$. For each $j=1,2$, we take a holomorphic function $h_{j}$, on a neighborhood $\mathcal{V}_{h_{j}}$ of $\Delta_{j}$, i.e. $h_{j} \in H\left(\mathcal{V}_{h_{j}}\right)$. Let us define the system of measures $\left(\sigma_{1}, \sigma_{2}\right)$ where $\sigma_{1}$ and $\sigma_{2}$ have the differential form

$$
d \sigma_{j}(x)=\frac{h_{j}(x) d x}{\sqrt{\left(x-c_{1, j}\right)\left(c_{2, j}-x\right)}}, \quad x \in \Delta_{j}, \quad j=1,2 .
$$

This system $\left(\sigma_{1}, \sigma_{2}\right)$ belongs to the class of Angelesco systems introduced by Angelesco in [1]. Fix a multi-index $\boldsymbol{n}=\left(n_{1}, n_{2}\right)$, we say that a polynomial $Q_{n} \not \equiv 0$ is a type II multiple-orthogonal polynomial corresponding to a system $\left(\sigma_{1}, \sigma_{2}\right)$, if $\operatorname{deg} Q_{n} \leq|\boldsymbol{n}|=n_{1}+n_{2}$ and $Q_{\boldsymbol{n}}$ satisfies the following orthogonality conditions

$$
\int_{\Delta_{j}} x^{\nu} Q_{n}(x) d \sigma_{j}(x)=0, \quad \nu=0, \ldots, n_{j}-1, \quad j=1,2 .
$$

It is well known that for any multi-index $\boldsymbol{n}=\left(n_{1}, n_{2}\right)$, the polynomial $Q_{\boldsymbol{n}}$ has exactly $n_{1}+n_{2}$ simple zeros lying in $\stackrel{\circ}{\Delta}_{1} \cup \stackrel{\circ}{\Delta}_{2}$, where $\stackrel{\circ}{\Delta}_{j}$ denotes the interior set of $\Delta_{j}, j=1,2$. Our propose in the present article consists in obtaining results about the strong asymptotic development of sequences of multi-orthogonal

[^0]polynomials $\left\{Q_{n}: \boldsymbol{n} \in \mathbb{Z}^{2}\right\}$. An effective method for such study with this kind of "very well" measures, is analyzing of the Riemann-Hilbert problem for multi-orthogonal polynomials, which was introduced in [5]. Let us consider a $3 \times 3$ square matrix, $Y$, whose entries are complex functions $Y_{s, k}: \mathbb{C} \rightarrow \mathbb{C}$, $s, k=1,2,3$. Given a point $x \in \stackrel{\circ}{\Delta}_{1} \cup \stackrel{\circ}{\Delta}_{2}$, the following matricial limits, where $z \in \mathbb{C} \backslash\left(\Delta_{1} \cup \Delta_{2}\right)$ tending to $x$, represent the formal pontual limits of all entries of $Y$ at the same time:
\[

$$
\begin{aligned}
& \lim _{z \rightarrow x} Y(z)=Y_{+}(x), \quad \Im m(z)>0 \\
& \lim _{z \rightarrow x} Y(z)=Y_{-}(x), \quad \Im m(z)<0
\end{aligned}
$$
\]

Let $\delta_{s, k}$ denote the Kroneker delta function. Let us look for a matrix function $Y$ which satisfies the following conditions:
(1) The entries of $Y, Y_{s, k}$, belongs to $H\left(\mathbb{C} \backslash\left(\Delta_{1} \cup \Delta_{2}\right)\right)$, which we write as $Y \in H\left(\mathbb{C} \backslash\left(\Delta_{1} \cup \Delta_{2}\right)\right)$;
(2) For each $\Delta_{j}, j=1,2$, the so called jump condition takes place

$$
Y_{+}(x)=Y_{-}(x)\left(\begin{array}{ccc}
1 & \frac{2 \pi i \delta_{1, j} h_{1}(x) d x}{\sqrt{\left(x-c_{1,1}\right)\left(c_{2,1}-x\right)}} & \frac{2 \pi i \delta_{2, j} h_{2}(x) d x}{\sqrt{\left(x-c_{1,2}\right)\left(c_{2,2}-x\right)}} \\
0 & 1 & 1
\end{array}\right), \quad x \in \stackrel{\circ}{\Delta}_{j} ;
$$

(3) Given a multi-index $\boldsymbol{n}=\left(n_{1}, n_{2}\right)$, we require the following asymptotic condition at infinity,

$$
Y(z)\left(\begin{array}{ccc}
z^{-|n|} & 0 & 0 \\
0 & z^{n_{1}} & 0 \\
0 & 0 & z^{n_{2}}
\end{array}\right)=\mathbb{I}+\mathcal{O}(1 / z) \quad \text { as } \quad z \rightarrow \infty
$$

where $\mathbb{I}$ is the identity matrix with rank 3 ;
(4) For each $i, j=1,2$, we set the following behavior around the endpoints $c_{i, j}$,

$$
Y(z)=\mathcal{O}\left(\begin{array}{lll}
1 & \delta_{2, j}+\frac{\delta_{1, j}}{\sqrt{\left|z-c_{i, j}\right|}} & \delta_{1, j}+\frac{\delta_{2, j}}{\sqrt{\left|z-c_{i, j}\right|}} \\
1 & \delta_{2, j}+\frac{\delta_{1, j}}{\sqrt{\left|z-c_{i, j}\right|}} & \delta_{1, j}+\frac{\delta_{\delta_{j, j}}}{\sqrt{\left|z-c_{i, j}\right|}} \\
1 & \delta_{2, j}+\frac{\delta_{1, j}}{\sqrt{\left|z-c_{i, j}\right|}} & \delta_{1, j}+\frac{\delta_{2, j}}{\sqrt{\left|z-c_{i, j}\right|}}
\end{array}\right) .
$$

This problem, which consists in finding the matrix function $Y$, was called in [5] a Riemann-Hilbert problem for type II multiple orthogonal polynomials, and
for the system of measures ( $\sigma_{1}, \sigma_{2}$ ), RHP in short. The solution $Y$ is unique and has the form

$$
Y(z)=\left(\begin{array}{ccc}
Q_{n}(z) & -\int_{\Delta_{1}} Q_{n}(x) \frac{d \sigma_{1}(x)}{z-x} & -\int_{\Delta_{2}} Q_{n} \frac{d \sigma_{2}(x)}{z-x}  \tag{1}\\
d_{1} Q_{n_{-}^{1}}(z) & -\int_{\Delta_{1}} Q_{n_{-}}(x) \frac{d \sigma_{1}(x)}{z-x} & -\int_{\Delta_{2}} Q_{n_{-}^{1}} \frac{d \sigma_{2}(x)}{z-x} \\
d_{2} Q_{n_{-}^{2}}^{2}(z) & -\int_{\Delta_{1}} Q_{n_{-}^{2}}(x) \frac{d \sigma_{1}(x)}{z-x} & -\int_{\Delta_{2}} Q_{n_{-}^{2}}^{2} \frac{d \sigma_{2}(x)}{z-x}
\end{array}\right)
$$

with ${d_{i}}^{-1}=-\int_{\Delta_{i}} x^{n_{i}-1} Q_{n_{-}^{i}}(x) d \sigma_{i}(x), i=1,2$, and if $\boldsymbol{n}=\left(n_{1}, n_{2}\right), \boldsymbol{n}_{-}^{1}=$ $\left(n_{1}-1, n_{2}\right)$ and $\boldsymbol{n}_{-}^{2}=\left(n_{1}, n_{2}-1\right)$.

The key of our procedure is based in finding the relationship between $Y$ and a matrix function $R$ which is the solution of another RHP with the following formulation. Suppose that $\gamma$ is a closed simple and smooth contour on the complex plane $\mathbb{C}$, then find a matrix function, $R$, such that:
(1) $R: \mathbb{C} \rightarrow \mathbb{C}^{3 \times 3}$ belongs to $H(\mathbb{C} \backslash \gamma)$;
(2) $R_{+}(\xi)=R_{-}(\xi) V_{n}(\xi), \xi \in \gamma$;
(3) $R(z) \rightarrow \mathbb{I}$ as $z \rightarrow \infty$,
where $V_{n}$ is a $3 \times 3$ matrix function, which is called the jump matrix.
Given an arbitrary $3 \times 3$ matrix function $K=\left[K_{s, k}\right]_{s, k}, s, k=1,2,3$, defined on a open set $\Omega \subset \mathbb{C}$ let us denote by $\|K\|$ (respectively, $\|K\|_{\Omega}$ ) the matrix infinity norm which consists in the maximum sum of row's entries modulus, defined for $3 \times 3$ matrices, i.e.

$$
\left.\|K\|=\max _{s=1,2,3} \sum_{k=1}^{3}\left|K_{s, k}\right|, \quad \text { (respectively, } \quad\|K\|_{\Omega}=\sup _{\Omega}\|K\|\right) .
$$

Theorem 1 (See Theorem 3.1 in [7]). Suppose that $\Omega$ is an open set containing $\gamma$. In condition (2) of the RHP for $R$, let us require for $V_{n} \in H(\Omega)$ that there exist constants $C$ and $\delta_{n}>0$ for which

$$
\left\|V_{n}-\mathbb{I}\right\|_{\Omega}<\delta_{n}
$$

Then, any solution of the RHP for $R$ satisfies that

$$
\|R(z)-\mathbb{I}\|<C\left\|V_{n}-\mathbb{I}\right\|_{\Omega} \quad \text { for every } \quad z \in \mathbb{C} \backslash \gamma
$$

Notice that if we know the relationship between $R$ and $Y$ and if we can also describe the development of $R$ when $|\boldsymbol{n}| \rightarrow \infty$, we would have a description for the development of all entries of $Y$ when $|\boldsymbol{n}| \rightarrow \infty$, particularly for $Y_{1,1}(z)=Q_{n}(z)$.

The RHP for $Y$ is not normalized in the sense that the conditions (3) at infinity for $Y$ and $R$ are different. In order to normalize the RHP, we are going to modify $Y$ in such a way that we set another RHP with the same contours (possibly different jump conditions), for which the solution tends to the identity matrix as $z \rightarrow \infty$. For normalizing we need to take into account the behavior of $Y(z)$ for large $z$. This behavior depends on the distribution of the zeros of the multiple-orthogonal polynomials. The zero distribution of the orthogonal polynomials is usually given by an extremal problem in logarithmic potential theory. In section 2 we introduce some concepts and results which we will need about this theory and we will normalized the Riemann-Hilbert problem at infinity. In section 3 such a Riemann-Hilbert problem with oscillatory and exponentially decreasing jumps can be analyzed by using the steepest descent method introduced by Deift and Zhou (see [3, 4]).

## 2. The equilibrium problem and the normalization at infinity

Let us fix $j \in\{1,2\} . \mathcal{M}_{1 / 2}\left(\Delta_{j}\right)$ denotes the set of all finite Borel measures whose supports, i.e. $\operatorname{supp}(\cdot)$, are contained in $\Delta_{j}$ with total variation $1 / 2$. Take $\mu_{j} \in \mathcal{M}_{1 / 2}\left(\Delta_{j}\right)$ and define its logarithmic potential as follows

$$
V^{\mu_{j}}(z)=\int \log \frac{1}{|z-x|} d \mu_{j}(x), \quad z \in \mathbb{C}
$$

For each pair of measures $\left(\mu_{1}, \mu_{2}\right)$, where $\mu_{j} \in \mathcal{M}_{1 / 2}\left(\Delta_{j}\right), j=1,2$, we define the quantities

$$
m_{j}\left(\mu_{1}, \mu_{2}\right)=\min _{x \in \Delta_{j}}\left(2 V^{\mu_{j}}(x)+V^{\mu_{k}}(x)\right), \quad j, k=1,2, \quad j \neq k
$$

The following Proposition is deduced immediately from the results of [6].
Proposition 1. There exists a unique pair $\left(\bar{\mu}_{1}, \bar{\mu}_{2}\right) \in \mathcal{M}_{1 / 2}\left(\Delta_{1}\right) \times \mathcal{M}_{1 / 2}\left(\Delta_{2}\right)$, which satisfies for $j, k=1,2$

$$
2 V^{\bar{\mu}_{j}}(x)+V^{\bar{\mu}_{k}}(x)=m_{j}\left(\bar{\mu}_{1}, \bar{\mu}_{2}\right)=m_{j}, \quad x \in \operatorname{supp}\left(\bar{\mu}_{j}\right)=\Delta_{j}, \quad j \neq k .
$$

For each $j=1,2$ the measure $\bar{\mu}_{j}$ has the following differential form

$$
d \bar{\mu}_{1}(x)=\frac{\rho_{1}(x) d x}{\sqrt{\left(x-c_{1,1}\right)\left(c_{2,1}-x\right)}}, \quad d \bar{\mu}_{2}(x)=\frac{\rho_{2}(x) d x}{\sqrt{\left(x-c_{1,2}\right)\left(c_{2,2}-x\right)}}
$$

where $\rho_{j} \in H\left(\mathcal{V}_{\rho_{j}}\right)$, with $\mathcal{V}_{\rho_{j}}$ denoting an open set which contains $\Delta_{j}$.

The pair $\left(\bar{\mu}_{1}, \bar{\mu}_{2}\right)$ is called extremal or equilibrium pair of measures with respect to $\left(\Delta_{1}, \Delta_{2}\right)$. Let us denote for each $j=1,2$ the analytic potentials

$$
g_{j}(z)=\int_{\Delta_{j}^{*}} \log (z-x) d \bar{\mu}_{j}(x)=-V^{\bar{\mu}_{j}}(z)+i \int_{\Delta_{j}^{*}} \arg (z-x) d \mu_{j}(x),
$$

where $\Delta_{j}^{*}$ is the support of the extremal measure, $\bar{\mu}_{j}$, that coincides in our case with $\Delta_{j}$, for $j=1,2$. Substituting the potential logarithmic in Proposition 1 we obtain for each $j, k=1,2$ with $j \neq k$ that

$$
-\left[g_{j+}+g_{j-}\right](x)-g_{k-}(x)=m_{j}, \quad x \in \Delta_{j} .
$$

Observe that

$$
\left(g_{j+}-g_{j-}\right)(x)=\left\{\begin{array}{lll}
0 & \text { if } \quad c_{2, j} \leq x \\
i \pi & \text { if } \quad c_{1, j} \geq x \\
2 i \pi \int_{x}^{c_{2, j}} d \bar{\mu}_{j}(t) & \text { if } \quad x \in \Delta_{j}
\end{array}\right.
$$

In what follows all the multi-indices will have the form $\boldsymbol{n}=(n, n)$. Let us introduce the matrices

$$
G(z)=\left(\begin{array}{ccc}
e^{-2 n\left(g_{1}(z)+g_{2}(z)\right)} & 0 & 0  \tag{2}\\
0 & e^{2 n g_{1}(z)} & 0 \\
0 & 0 & e^{2 n g_{2}(z)}
\end{array}\right), \quad L=\left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & e^{-2 n m_{1}} & 0 \\
0 & 0 & e^{-2 n m_{2}}
\end{array}\right) .
$$

We define the matrix function $T=L Y G L^{-1}$, where $L, G$ are as in (2) and $Y$ is given by (1). Hence $T$ is the unique solution of the RHP:
(1) $T \in H\left(\mathbb{C} \backslash\left(\stackrel{\circ}{\Delta}_{1} \cup \stackrel{\circ}{\Delta}_{2}\right)\right)$;
(2) $T_{+}(x)=T_{-}(x) M(x), x \in \Delta_{1} \cup \Delta_{2}$;
(3) $T(z)=\mathbb{I}+\mathcal{O}(1 / z)$ as $z \rightarrow \infty$;
(4) $T$ and $Y$ have the same behavior on the endpoints of the intervals $\Delta_{j}$, for $j=1,2$,
where the jump matrix $M$ has the form

$$
M(x)=\left(\begin{array}{ccc}
e^{-2 n i \pi \int_{x}^{c_{2, j}} d \bar{\mu}_{j}(t)} & 2 \delta_{j, 1} \pi i w_{1}(x) & 2 \delta_{j, 2} \pi i w_{2}(x) \\
0 & e^{2 n \delta_{j, i} i \pi \int_{x}^{c_{2,1}} d \bar{\mu}_{1}(t)} & 0 \\
0 & 0 & e^{2 n \delta_{j, 2} i \pi \int_{x}^{c_{2,2}} d \bar{\mu}_{2}(t)}
\end{array}\right), \quad x \in \stackrel{\circ}{\Delta}_{j} .
$$

## 3. The opening of the lens

For each $j=1,2$, let $\phi_{j}$ denote the function defined by

$$
\phi_{j}(z)=i \pi \int_{z}^{c_{2, j}} d \bar{\mu}_{j}(t) \quad \text { for } \quad z \in \mathcal{V}_{j}=\mathcal{V}_{\rho_{j}}\left(\Delta_{j}^{*}\right) \cap \mathcal{V}_{h_{j}}\left(\Delta_{j}\right)
$$

Notice that $\phi_{j+}(x)=i \pi \int_{x}^{c_{2, j}} d \bar{\mu}_{j}(t)$ is purely imaginary and its derivative

$$
\phi_{j+}^{\prime}(x)=-i \pi \frac{\rho_{j}(x)}{\sqrt{\left(x-c_{1, j}\right)\left(c_{2, j}-x\right)}},
$$

where $-\pi \rho_{j}(x) / \sqrt{\left(x-c_{1, j}\right)\left(c_{2, j}-x\right)}<0, \quad x \in \Delta_{j}$.
Rewrite $\phi_{j}(z)=U_{j}(z)+i V_{j}(z) \in H\left(\mathcal{V}_{j}\right)$. By the Cauchy-Riemann conditions we have that the real part of $\phi_{j}$, $\Re e \phi_{j}$, is an increasing function on any point $z \in \mathcal{V}_{j}$ with $\Im m(z)>0$. Since $\Re e \phi_{j}$ is zero in $\Delta_{j}$, it is positive in such point. Notice $\phi_{j+}(x)=-\phi_{j-}(x), x \in \stackrel{\circ}{\Delta}_{j}$, hence we can proceed analogously when $\Im m(z)<0$.
We analyze the jump function in $\Delta_{j}$, i.e.

$$
\begin{aligned}
& M(x)=\left(\begin{array}{ccc}
e^{-2 n \phi_{j+}(x)} & 2 \pi i \delta_{j, 1} w_{1}(x) & 2 \pi i \delta_{j, 2} w_{2}(x) \\
0 & e^{-2 \delta_{j, 1} n \phi_{1}(x)} & 0 \\
0 & 0 & e^{-2 \delta_{j, 2} n \phi_{2-}(x)}
\end{array}\right)=\left(\begin{array}{ccc}
1 & 0 & 0 \\
-\frac{\delta_{j, 1} e^{-2 n \phi_{1}-(x)}}{2 \pi i w_{1}(x)} & 1 & 0 \\
-\frac{\delta_{j, 2} e^{-2 n n \phi_{2}-(x)}}{2 \pi i w_{2}(x)} & 0 & 1
\end{array}\right) \\
& \times\left(\begin{array}{ccc}
0 & 2 \pi i \delta_{j, 1} w_{1}(x) & 2 \pi i \delta_{j, 2} w_{2}(x) \\
-\frac{\delta_{1, j}}{2 \pi w_{1}(x)} & \delta_{j, 2} & 0 \\
-\frac{\delta_{2, j}}{2 \pi i w_{2}(x)} & 0 & \delta_{j, 1}
\end{array}\right)\left(\begin{array}{ccc}
1 & 0 & 0 \\
-\frac{\delta_{j, 1} e^{-2 n \phi_{1}(x)}}{2 \pi i w_{1}(x)} & 1 & 0 \\
-\frac{\delta_{j, 2} e^{-2 n \phi_{2}+(x)}}{2 \pi i w_{2}(x)} & 0 & 1
\end{array}\right) .
\end{aligned}
$$

Now we are going to follow an analogous procedure as in section 9 of [7]. Let us fix a $\delta>0$ such that the intervals $\left[c_{1, j}-\delta, c_{1, j}\right]$ and $\left[c_{2, j}, c_{2, j}+\delta\right]$ are subsets of $\mathcal{V}_{j}, j=1,2$, and for each interval let us define two curves $\Sigma_{j+}, \Sigma_{j-}$ in $\mathcal{V}_{j}$, which goes from $c_{1, j}$ to $c_{2, j}$, where $\Sigma_{j \pm}=\left[c_{1, j}-\delta, c_{1, j}\right] \cup \Sigma_{j \pm}^{*} \cup\left[c_{2, j}, c_{2, j}+\delta\right]$, with the elements of the curves $\Sigma_{j \pm}^{*}$ satisfying that if $z \in \Sigma_{j \pm}^{*}$, then $0< \pm \Im m(z)$ (cf. Figure 1). Set $\Gamma_{j \pm}$ the domains that lie between $\Sigma_{j \pm}$ and $\Delta_{j}$. Let us introduce the matrix function $S$, defined by

$$
S(z)=T\left(\begin{array}{ccc}
1 & 0 & 0  \tag{3}\\
\mp \frac{\delta_{1, j} e^{-2 n \phi_{1}-(z)}}{2 \pi i w_{1}(z)} & 1 & 0 \\
\mp \frac{\delta_{2, j} e^{-2 n \phi_{2}-(z)}}{2 \pi i w_{2}(z)} & 0 & 1
\end{array}\right), \quad z \in \Gamma_{j \pm}, \quad \text { and } S \equiv T, \quad \text { outside. }
$$

Hence on open intervals $] c_{1, j}-\delta, c_{1, j}[$ and $] c_{2, j}, c_{2, j}+\delta[$ there are two combined jumps. For the function $w_{j}$ we have that

$$
w_{j+}(x)=-w_{j-}(x), \quad x \in \mathbb{R} \cap \mathcal{V}\left(\Delta_{j}\right) \backslash \Delta_{j}, \quad j=1,2
$$

Observe that

$$
\left(\begin{array}{ccc}
1 & 0 & 0 \\
\frac{\delta_{j, 1} e^{-2 n \phi_{1}}}{2 \pi i w_{1}} & 1 & 0 \\
\frac{\delta_{j, 2} e^{-2 n \phi_{2}}}{2 \pi i w_{2-}} & 0 & 1
\end{array}\right)\left(\begin{array}{ccc}
1 & 0 & 0 \\
\frac{\delta_{j, 1} e^{-2 n \phi_{1}}}{2 \pi i w_{1}+} & 1 & 0 \\
\frac{\delta_{j, 2} e^{2 n \phi_{2}}}{2 \pi i w_{2+}} & 0 & 1
\end{array}\right)=\left(\begin{array}{ccc}
1 & 0 & 0 \\
\left(\frac{1}{w_{1-}}+\frac{1}{w_{1+}}\right. \\
\left(\frac{1}{w_{2-}}+\frac{1}{w_{2+}}\right) \frac{\delta_{1, j} e^{-2 n \phi_{1}}}{2 \pi i} & 1 & 0 \\
\frac{\delta_{2, j} e^{-2 n \phi_{1}}}{2 \pi i} & 0 & 1
\end{array}\right)=\mathbb{I},
$$

which means that $S$, defined by (3), is analytic function across $] c_{1, j}-\delta, c_{1, j}[$ and $] c_{2, j}, c_{2, j}+\delta\left[, j=1,2\right.$. Let $\gamma_{j}, j=1,2$, be closed contours with the clockwise direction, such that for each $j=1,2, \gamma_{j}=\Sigma_{j-}^{*} \cup \Sigma_{j+}^{*}$. We have changed the direction of the curve $\Sigma_{j+}^{*}$. The function $S$ satisfies the RHP:
(1) $S \in H\left(\mathbb{C} \backslash \cup_{j=1,2}\left(\Delta_{j} \cup \gamma_{j}\right)\right)$;
(2) The jump conditions for $j=1,2$ are,

$$
\begin{aligned}
& S_{+}(x)=S_{-}(x)\left(\begin{array}{ccc}
0 & 2 \delta_{1, j} \pi i w_{1}(x) & 2 \delta_{2, j} \pi i w_{2}(x) \\
-\frac{\delta_{1, j}}{2 \pi i w_{1}(x)} & \delta_{2, j} & 0 \\
-\frac{\delta_{2, j}}{2 \pi i w_{2}(x)} & 0 & \delta_{1, j}
\end{array}\right) \text { if } x \in \stackrel{\circ}{\Delta}_{j}, \\
& S_{+}(z)=S_{-}(z)\left(\begin{array}{ccc}
1 & 0 & 0 \\
\frac{ \pm \delta_{1, j} e^{-2 n \phi_{1}(z)}}{2 \pi i_{1}(z)} & 1 & 0 \\
\frac{ \pm \delta_{2, j} e^{2}-2 \phi_{2}(z)}{2 \pi i w_{2}(z)} & 0 & 1
\end{array}\right) \text { if } z \in \gamma_{j} \cap\{ \pm \Im m z<0\} ;
\end{aligned}
$$

(3) $S(z)=\mathbb{I}+\mathcal{O}(1 / z)$ as $z \rightarrow \infty$;
(4) The conditions for the endpoints are the same as for $T$.


Figure 1. Opening of lens

Now, we consider the limiting problem, because for the matrix $S$ the jump matrix function on each $\gamma_{j}$ for $j=1,2$ tends to the identity matrix when $|n| \rightarrow \infty$. We look for the matrix function $N$ which satisfies the following RHP:
(1) $N \in H\left(\mathbb{C} \backslash\left(\Delta_{1} \cup \Delta_{2}\right)\right)$;
(2) The jump conditions in $\grave{\Delta}_{j}$ for $j=1,2$ are,

$$
N_{+}(x)=N_{-}(x)\left(\begin{array}{ccc}
0 & 2 \delta_{1, j} \pi i w_{1}(x) & 2 \delta_{2, j} \pi i w_{2}(x)  \tag{4}\\
-\frac{\delta_{1, j}}{2 \pi u_{2 j}(x)} & \delta_{2, j} & 0 \\
-\frac{\delta_{2}}{2 \pi i w_{2}(x)} & 0 & \delta_{1, j}
\end{array}\right) ;
$$

(3) $N(z)=\mathbb{I}+\mathcal{O}(1 / z)$ as $z \rightarrow \infty$;
(4) $N$ satisfies the same conditions for the endpoints as $S$.

Set $N=K D$, where $D$ is a diagonal matrix function and $K=\left[K_{k, l}\right]_{k, l}$, $k, l=1,2$, is the solution of the RHP:
(1) $K \in H\left(\mathbb{C} \backslash\left(\Delta_{1} \cup \Delta_{2}\right)\right)$;
(2) The jump conditions in $\grave{\Delta}_{j}$ for $j=1,2$ are, because of (4),

$$
K_{+}(x)=K_{-}(x)\left(\begin{array}{ccc}
0 & \frac{2 \delta_{1, j} \pi i}{\sqrt{\left(c_{2,1}, x\right)\left(x-c_{1,1}\right)}} & \frac{2 \delta_{2, j} \pi i}{\sqrt{\left(c_{2,2}-x\right)\left(x-c_{1,2}\right)}}  \tag{5}\\
\frac{-\delta_{1, j} \sqrt{\left(c_{2,2}-x\right)\left(x-c_{1,1)}\right)}}{\left(\delta_{2, j} \sqrt{\left(c_{2,2}-x\right)\left(x-c_{1,2}\right)}\right.} & \delta_{2, j} & 0 \\
\frac{-\delta_{i}}{2 \pi i} & \delta_{1, j}
\end{array}\right) ;
$$

(3) $K(z)=\mathbb{I}+\mathcal{O}(1 / z)$ as $z \rightarrow \infty$;
(4) $K$ and $N$ have the same conditions for the endpoints.

Analogously to the ideas in [2], let us choose the branches of the square root which glue along the intervals $\Delta_{j}, j=1,2$, i.e.

$$
\left(\sqrt{\left(x-c_{2, j}\right)\left(x-c_{1, j}\right)}\right)_{+}=-\left(\sqrt{\left(x-c_{2, j}\right)\left(x-c_{1, j}\right)}\right)_{-}, x \in{\stackrel{\circ}{\Delta_{j}}}_{j}, j=1,2 .
$$

For each $i=1,2,3$, we rewrite (5) as

$$
\left\{\begin{array}{l}
\left(\frac{\sqrt{\left(z-c_{2,1}\right)\left(z-c_{1,1}\right)}}{2 \pi} K_{i, 2}(z)\right)_{ \pm}(x)=\left(K_{i, 1}(z)\right)_{\mp}(x) \\
\left(K_{i, 3}\right)_{+}(x)=\left(K_{i, 3}\right)_{-}(x)
\end{array} \quad, \quad x \in \stackrel{\circ}{\Delta}_{1}\right.
$$

$$
\left\{\begin{array}{l}
\left(\frac{\sqrt{\left(z-c_{2,2}\right)\left(z-c_{1,2}\right)}}{2 \pi} K_{i, 3}(z)\right)_{ \pm}(x)=\left(K_{i, 1}(z)\right)_{\mp}(x) \quad, \quad x \in \stackrel{\circ}{\Delta}_{2} \\
\left(K_{i, 2}\right)_{+}(x)=\left(K_{i, 2}\right)_{-}(x)
\end{array}\right.
$$

and we denote

$$
\psi_{0}^{i}(z)=K_{i, 1}(z), \quad \text { and } \quad \psi_{j}^{i}(z)=\frac{\sqrt{\left(z-c_{2, j}\right)\left(z-c_{1, j}\right)}}{2 \pi} K_{i, j+1}(z), j=1,2
$$

Then from the relations (5), we may interpret each row $i=1,2,3$ of such matrix $K$ as a function defined on a Riemann surface. Let $\mathcal{R}$ define the Riemann surface which has two cuts. One of them connects the two branch points $c_{1,1}$ and $c_{2,1}$ with the cut in the interval $\Delta_{1}=\left[c_{1,1}, c_{2,1}\right]$. The other cut is made in the interval $\Delta_{2}=\left[c_{1,2}, c_{2,2}\right]$, to connect the two other branch points $c_{1,2}$ and $c_{2,2}$. The sheet $\mathcal{R}_{0}$ is glued to another sheet $\mathcal{R}_{1}$ along the cut $\Delta_{1}$, and $\mathcal{R}_{0}$ is also glued to $\mathcal{R}_{2}$ along the interval $\Delta_{2}$. Let us denote by $\psi^{i}, i=1,2,3$, three multi-valued function $\psi^{i}=\left(\psi_{0}^{i}, \psi_{1}^{i}, \psi_{2}^{i}\right)$, such that its components $\psi_{l}^{i}, i=1,2,3, l=0,1,2$, map the corresponding sheet $\mathcal{R}_{l}$ on $\mathbb{C}$, and satisfy for $j=1,2$

$$
\begin{equation*}
\psi_{0 \pm}^{i}(x)=\psi_{j \mp}^{i}(x), \quad x \in \stackrel{\circ}{\Delta}_{j}, \quad \psi_{j}^{i}(z)=\mathcal{O}(1) \quad \text { as } \quad z \rightarrow c_{k, j}, \quad k=1,2 ; \tag{6}
\end{equation*}
$$

for $j, k=1,2, \psi_{0}^{i}(x)=\mathcal{O}(1)$, as $z \rightarrow c_{k, j}$; around the infinity, and $j=1,2$,

$$
\psi_{0}^{i}(z)=\delta_{i, 1}+\mathcal{O}(1 / z), \quad \psi_{j}^{i}(z)=z \delta_{i, j+1}+\mathcal{O}(1) \quad \text { as } \quad z \rightarrow \infty
$$

The equalities (6) are equivalent to

$$
\left(\psi_{0}^{i} \psi_{j}^{i}\right)_{+}(x)=\left(\psi_{0}^{i} \psi_{j}^{i}\right)_{-}(x), \quad x \in \stackrel{\circ}{\Delta}_{j}, \quad j=1,2
$$

From Liouville's theorem, it is easy to see that $\left(\psi_{0}^{i} \psi_{1}^{i} \psi_{2}^{i}\right)(z) \equiv 1, z \in \overline{\mathbb{C}}$. This implies that $\psi_{l}^{i}, l=0,1,2, i=1,2,3$, do not become zero. Hence

$$
\left(\psi_{0}^{i} \psi_{k}^{i}\right)(z)=\frac{1}{\psi_{l}^{i}(z)} \in H\left(\overline{\mathbb{C}} \backslash \Delta_{l}\right), \quad l, k=1,2, \quad l \neq k, \quad i=1,2,3
$$

We obtain that for each $l=1,2$, that is $x \in \stackrel{\circ}{\Delta}_{l}$,

$$
\psi_{0+}^{i}(x) \psi_{k}^{i}(x)=\frac{1}{\psi_{l+}^{i}(z)} \quad \text { or equivalently } \quad\left(\psi_{l+}^{i} \psi_{l-}^{i}\right)(x)=\frac{1}{\psi_{k}^{i}(x)}
$$

That yields to the problems for $\psi_{l}^{i}, l=1,2$ :

- $\psi_{l}^{i} \in H\left(\overline{\mathbb{C}} \backslash \Delta_{l}\right)$;
- $\left(\psi_{l+}^{i} \psi_{l-}^{i}\right)(x)=1 / \psi_{k}^{i}(x), \quad x \in \stackrel{\circ}{\Delta}_{l}$;
- $\psi_{l}^{i}(z)=z \delta_{i, l+1}+\mathcal{O}(1)$ as $z \rightarrow \infty$;
- $\psi_{l}^{i}(z)=\mathcal{O}(1)$ as $z \rightarrow c_{k, l}, k=1,2$.

This problem is equivalent to the system of integral equations:

$$
\begin{aligned}
& \psi_{0}^{i}(z)=\frac{1}{\psi_{1}^{i}(z) \psi_{2}^{i}(z)} \\
& \psi_{l}^{i}(z) \\
= & \exp \left(\frac{\sqrt{\left(z-c_{1, l}\right)\left(z-c_{2, l}\right)}}{2 \pi} \int_{\Delta_{l}} \frac{\log \psi_{k}^{i}(x)}{\sqrt{\left(x-c_{1, l}\right)\left(c_{2, l}-x\right)}} \frac{d x}{z-x}+\delta_{i, l+1} g_{\Delta_{l}}(z)\right)
\end{aligned}
$$

for $l=1,2$, with $z \in \overline{\mathbb{C}} \backslash \stackrel{\circ}{\Delta}_{l}$, and $g_{\Delta_{j}}$ is the analytic function which tends to $\infty$ as $\log z$, and whose real part vanishes in $\Delta_{j}, j=1,2$

Let us find the diagonal $3 \times 3$ matrix function $D$ with diagonal elements $D_{0}, D_{1}, D_{2}$, such that $N(z) \equiv K(z) D(z)$. The conditions (5) yield that entries of $D$ must satisfy the following conditions

$$
\left\{\begin{array}{l}
h_{j}(x) D_{0 \pm}(x)=D_{j \mp}(x) \\
D_{k+}(x)=D_{k-}(x)
\end{array} \quad \text { when } \quad x \in \stackrel{\circ}{\Delta}_{j}, \quad j, k=1,2, \quad k \neq j\right.
$$

Analogously to the function $\psi_{l}^{i}$, we obtain the following problem for the entries of $D$ :
i) $\left(D_{0} D_{1} D_{2}\right) \equiv 1$, which implies that for each $l=0,1,2, D_{l}$ does not become zero;
ii) $D_{l} \in H\left(\overline{\mathbb{C}} \backslash \Delta_{l}\right), l=1,2$;
iii) $\left(D_{l+} D_{l-}\right)(x)=h_{s}(x) /\left(D_{s}(x)\right), \quad s=1,2, \quad s \neq l, \quad x \in \stackrel{\circ}{\Delta}_{l}, l=1,2$;
iv) $D_{l}(z)=1+\mathcal{O}(1 / z)$ as $z \rightarrow \infty, l=1,2$,
v) $D_{l}(z)=\mathcal{O}(1)$ as $z \rightarrow c_{k, l}, k=1,2, l=1,2$.

This problem is equivalent to the following system of integral equations where $l=1,2$,

$$
\begin{aligned}
D_{l}(z) & =\exp \left(\frac{\sqrt{\left(z-c_{1, l}\right)\left(z-c_{2, l}\right)}}{2 \pi} \int_{\Delta_{l}} \frac{\log \left(\frac{D_{k}(x)}{h_{k}(x)}\right)}{\sqrt{\left(x-c_{1, l}\right)\left(c_{2, l}-x\right)}} \frac{d x}{z-x}\right) \\
D_{0}(z) & =\frac{1}{D_{1}(z) D_{2}(z)} \quad \text { with } \quad z \in \overline{\mathbb{C}} \backslash \stackrel{\circ}{\Delta}_{l}
\end{aligned}
$$

Let us take the three multi-valued function $\left(D_{0}(z), D_{1}(z), D_{2}(z)\right)$. Notice that for each $l=0,1,2$, the function $D_{l}$ is another function which maps the sheet $\mathcal{R}_{l}$ on $\mathbb{C}$. In our case the components of functions $\left(D_{0}(z), D_{1}(z), D_{2}(z)\right)$ satisfy the conditions iv) and v) required for $D_{l}, l=0,1,2$. Finally the matrix function $N$ has the form

$$
N(z)=\left(\begin{array}{lll}
\left(D_{0} \psi_{0}^{1}\right)(z) & \frac{\left(D_{1} \psi_{1}^{1}\right)(z)}{\sqrt{\left(x-c_{1,1}\right)\left(c_{2,1}-x\right)}} & \frac{\left(D_{2} \psi_{2}^{1}\right)(z)}{\sqrt{\left(x-c_{1,2}\right)\left(c_{2,2}-x\right)}}  \tag{7}\\
\left(D_{0} \psi_{0}^{2}\right)(z) & \frac{\left(D_{1} \psi_{1}^{2}\right)(z)}{\sqrt{\left(x-c_{1,1}\right)\left(c_{2,1}-x\right)}} & \frac{\left(D_{2} \psi_{2}^{2}\right)(z)}{\sqrt{\left(x-c_{1,2}\right)\left(c_{2,2}-x\right)}} \\
\left(D_{0} \psi_{0}^{3}\right)(z) & \frac{\left(D_{1} \psi_{1}^{3}\right)(z)}{\sqrt{\left(x-c_{1,1}\right)\left(c_{2,1}-x\right)}} & \frac{\left(D_{2} \psi_{2}^{2}\right)(z)}{\sqrt{\left(x-c_{1,2}\right)\left(c_{2,2}-x\right)}}
\end{array}\right) .
$$

We define $R(z)=S(z) N^{-1}(z)$. Since $S$ and $N$ have the same jump across $\stackrel{\circ}{\Delta}_{j}$, $j=1,2$, hence $R_{+}(x)=R_{-}(x)$ for $x \in \stackrel{\circ}{\Delta}_{j}, j=1,2$. From the definition of $R$, and the endpoint conditions for $N$, we can also deduce that $c_{i, k}, i, k=1,2$, are a removable singularity. Hence $R$ is an analytic function across the full intervals $\Delta_{1}$ and $\Delta_{2}$, and it has jumps on the curve $\gamma$. Then we have the following RHP for $R$ :
(1) $R \in H\left(\mathbb{C} \backslash\left(\gamma_{1} \cup \gamma_{2}\right)\right)$;
(2) The jump conditions are for $j=1,2$

$$
R_{+}(z)=R_{-}(z)\left(\begin{array}{ccc}
1 & 0 & 0 \\
\pm \frac{\delta_{j, 1} e^{-2 n_{1} \phi_{1}(z)}}{2 \pi i i_{1}(z)} & 1 & 0 \\
\pm \frac{\delta_{j, 2} e^{-2 n p_{2} \phi_{2}(z)}}{2 \pi i w_{2}(z)} & 0 & 1
\end{array}\right) \quad \text { if } \quad z \in \gamma_{j} \cap\{ \pm \Im m z<0\}
$$

(3) $R(z)=\mathbb{I}+\mathcal{O}(1 / z)$.

Observe that matrix function $R$ satisfies the hypothesis of the Theorem 1. Then uniformly for $z \in \mathbb{C} \backslash\left(\gamma_{1} \cup \gamma_{2}\right)$, we have that $R(z)=\mathbb{I}+\mathcal{O}\left(e^{c|n|}\right)$, with $c>0$ and

$$
\begin{aligned}
Y(z)= & \left(\begin{array}{ccc}
1 & 0 & 0 \\
0 & e^{|\boldsymbol{n}| m_{1}} & 0 \\
0 & 0 & e^{|\boldsymbol{n}| m_{2}}
\end{array}\right)\left(\mathbb{I}+\mathcal{O}\left(e^{-c|\boldsymbol{n}|}\right)\right) \\
& \times N(z)\left(\begin{array}{ccc}
e^{|n|\left(g_{1}+g_{2}\right)}(z) & 0 & 0 \\
0 & e^{-|\boldsymbol{n}|\left(m_{1}+g_{1}(z)\right)} & 0 \\
0 & 0 & e^{-|n|\left(m_{2}+g_{2}\right)(z)}
\end{array}\right),
\end{aligned}
$$

where $N$ is given by (7).
Finally, we state the main result of this paper.

## Theorem 2.

$$
Y_{1,1}(z)=Q_{\boldsymbol{n}}(z)=D_{0}(z) \psi_{0}^{1}(z) e^{|\boldsymbol{n}|\left(g_{1}+g_{2}\right)(z)}\left(1+\mathcal{O}\left(e^{-c|\boldsymbol{n}|}\right)\right)
$$

as $|\boldsymbol{n}| \rightarrow \infty$.

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